

# Diffraction overview

Javier López-García

# Diffraction overview

- Diffraction history
- Crystallography
- Single-crystal diffraction
- Powder diffraction
- Magnetic diffraction
- Instruments

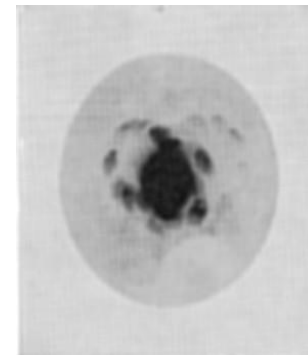
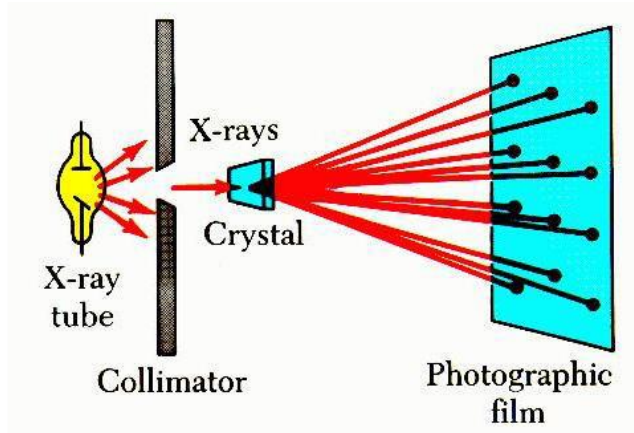
# Diffraction history



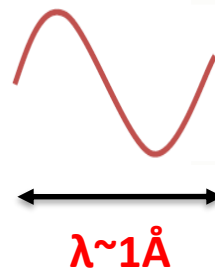
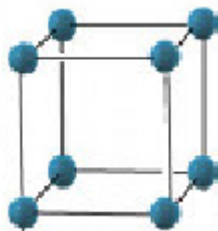
*Max von Laue*

Max von Laue  
(1879-1960)

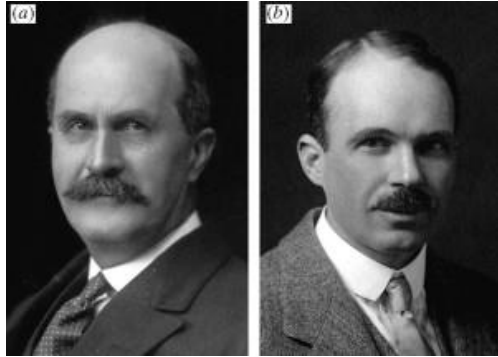
Nobel prize in  
1914 “for his  
discovery of  
the diffraction  
of X-Ray by  
crystals”



First diffractogram  
with X-Rays

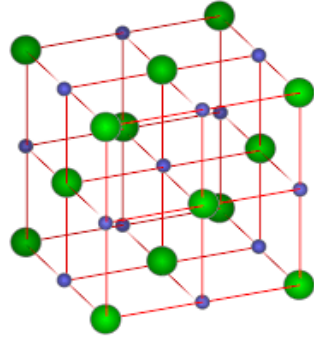


## Bragg family

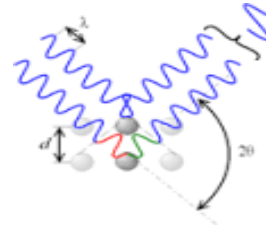


- a) William (1862-1942)
- b) Lawrence (1890-1971)

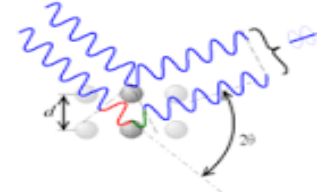
Nobel prize in 1915 “for their services in the analysis of crystal structure by means of X-rays”



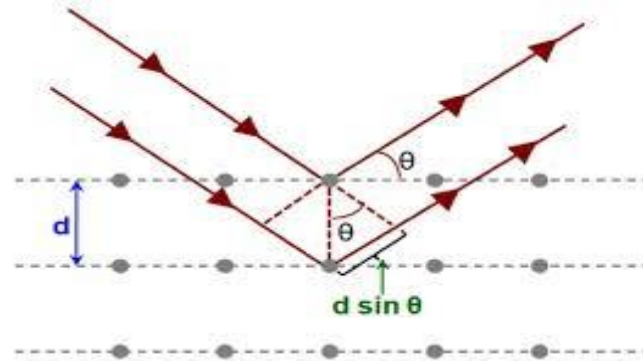
NaCl Fcc



Constructive interference



Destructive interference



Bragg law

$$2d_{hkl} \sin \theta_{hkl} = n\lambda$$



James Chadwick  
(1897-1974)

Discover the **neutron** in 1932

In 1936 it proves that **neutron** can be used to make **diffraction** due to wave-particle duality

Problem: weak source

Solution: nuclear reactors in 1942

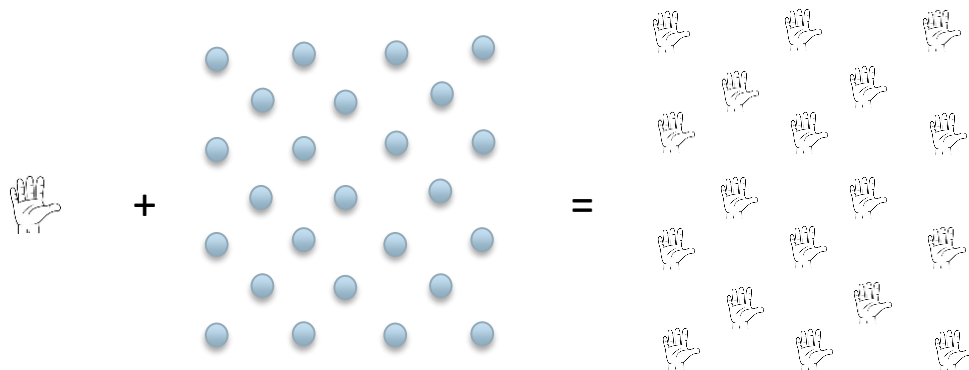
Fist neutron diffraction experiment was carried out by Clifford Shull and Ernest Wollan in 1946



Clifford Shull (right) and Ernest Wollan working with a double-crystal neutron spectrometer in 1949. Nobel prize in 1994 “for the development of the neutron scattering technique”

# Crystallography

What is a crystal?



Motif

Lattice

Crystal

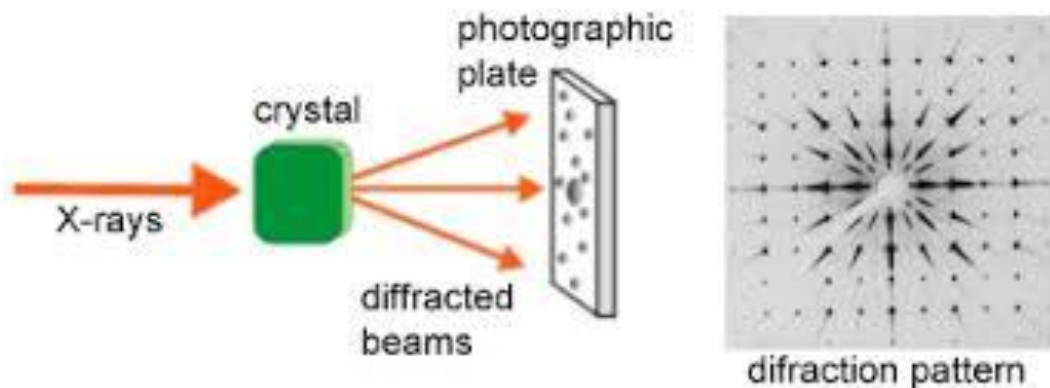
Crystal lattice structures

<b>CUBIC</b> $a = b = c$ $\alpha = \beta = \gamma = 90^\circ$	<b>P</b>	<b>I</b>	<b>F</b>
<b>TETRAGONAL</b> $a = b \neq c$ $\alpha = \beta = \gamma = 90^\circ$	<b>P</b>	<b>I</b>	
<b>ORTHORHOMBIC</b> $a \neq b \neq c$ $\alpha = \beta = \gamma = 90^\circ$	<b>P</b>	<b>I</b>	<b>F</b>
<b>HEXAGONAL</b> $a = b \neq c$ $\alpha = \beta = 90^\circ$ $\gamma = 120^\circ$	<b>P</b>		
<b>MONOCLINIC</b> $a \neq b \neq c$ $\alpha = \gamma = 90^\circ$ $\beta \neq 90^\circ$	<b>P</b>	<b>C</b>	
<b>TRICLINIC</b> $a \neq b \neq c$ $\alpha \neq \beta \neq \gamma \neq 90^\circ$	<b>P</b>		

**4 Types of Unit Cell**  
**P** = Primitive  
**I** = Body-Centred  
**F** = Face-Centred  
**C** = Side-Centred  
 +  
**7 Crystal Classes**  
 → **14 Bravais Lattices**

# Reciprocal space

The diffraction pattern obtained in the experiments is in **reciprocal space**



Reciprocal lattice is a network of points in the **Fourier space** defined for the vector:

$$\vec{\tau} = h\vec{a}^* + k\vec{b}^* + l\vec{c}^*$$

$$\vec{a}^* \cdot \vec{a} = 1$$

$$\vec{b}^* \cdot \vec{b} = 1$$

$$\vec{c}^* \cdot \vec{c} = 1$$

# Structure factor



For X-rays

$$F_{hkl} = \sum_j f_j e^{2\pi i(hx_j + ky_j + lz_j)} e^{-B_j(\sin^2\theta)/\lambda^2}$$

For neutron

$$F_{hkl} = \sum_j b_j e^{2\pi i(hx_j + ky_j + lz_j)} e^{-B_j(\sin^2\theta)/\lambda^2}$$

$F_{hkl}$ : structure factor

$f_j$ : X-rays form factor

$b_j$ : neutron form factor

hkl: Miller indices

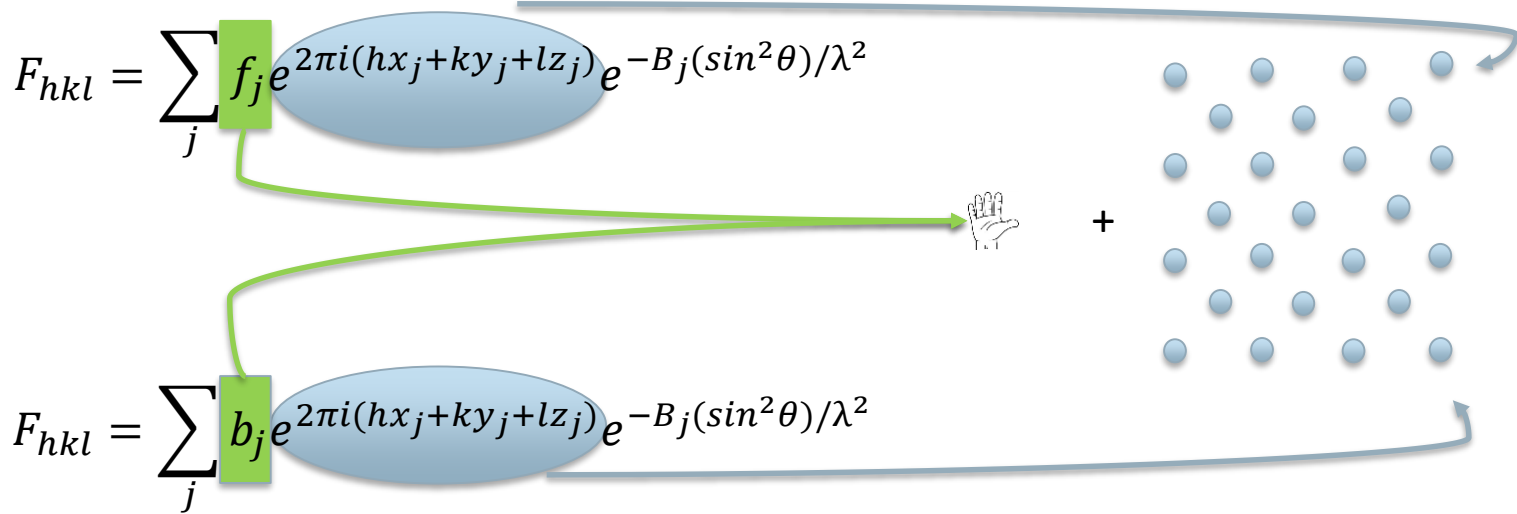
$x_j, y_j, z_j$ : coordinates of atom j

$B_j$ : thermal parameter

$\lambda$ : wavelength



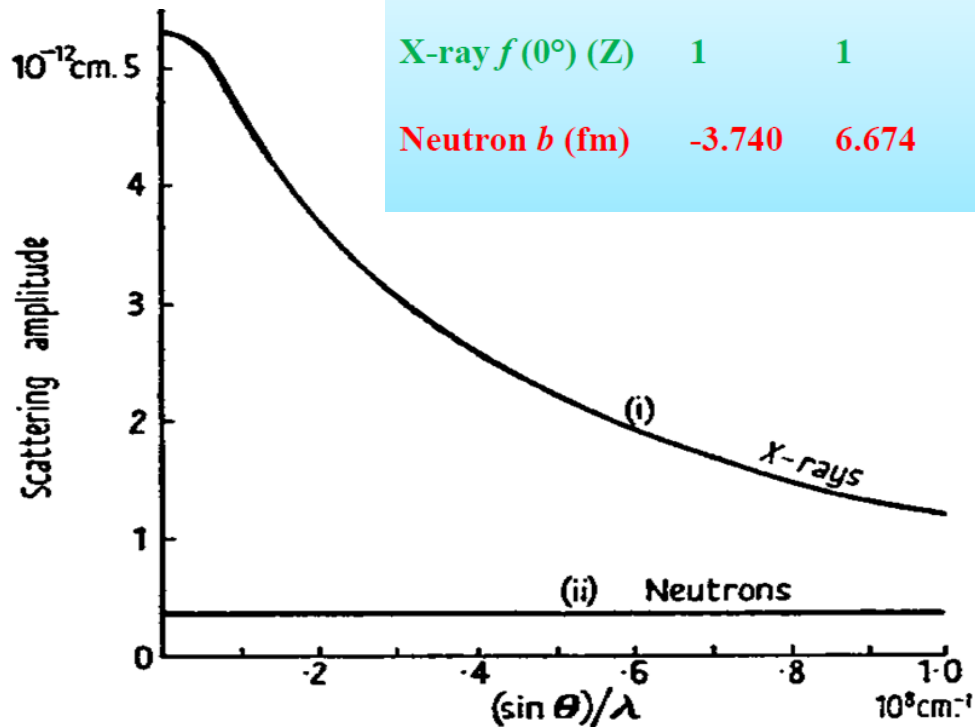
# Structure factor



# Structure factor



Element	H	D	U	Fe	Co	O	V	Ti
X-ray $f(0^\circ)$ (Z)	1	1	92	26	27	8	23	22
Neutron $b$ (fm)	-3.740	6.674	8.420	9.450	2.780	5.805	-0.3824	-3.438



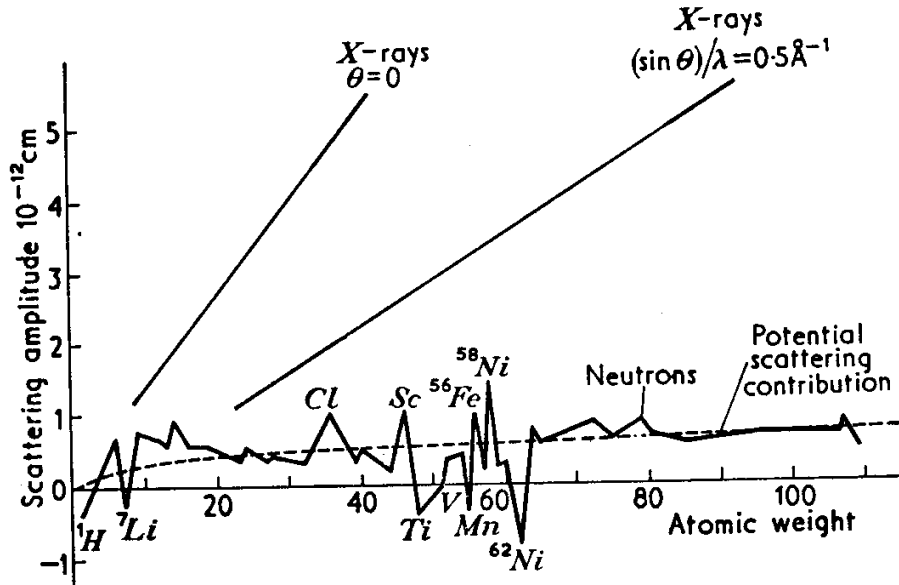
For same atom:

- X-Rays: scattering amplitude ↓ when  $(\sin \theta)/\lambda$  ↑
- Neutron: scattering amplitude is constant

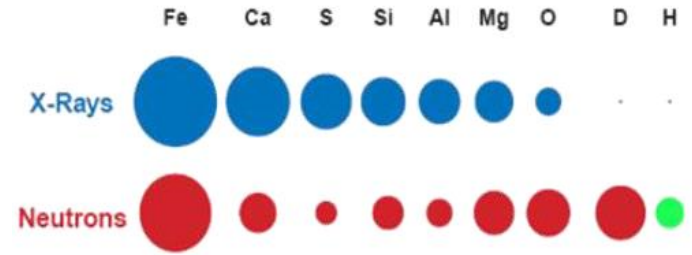
# Structure factor

In X-ray the magnitude of **f** is proportional to **Z**

In neutrons nuclear factor determine b



## Relative scattering factors

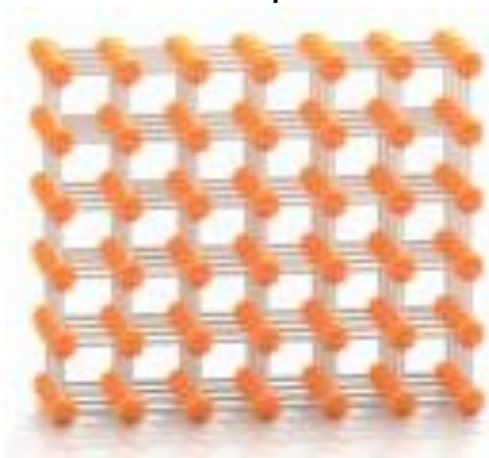


# Single-crystal diffraction

Crystal lattice **continuum** in the whole sample

All **oriented in the same direction**

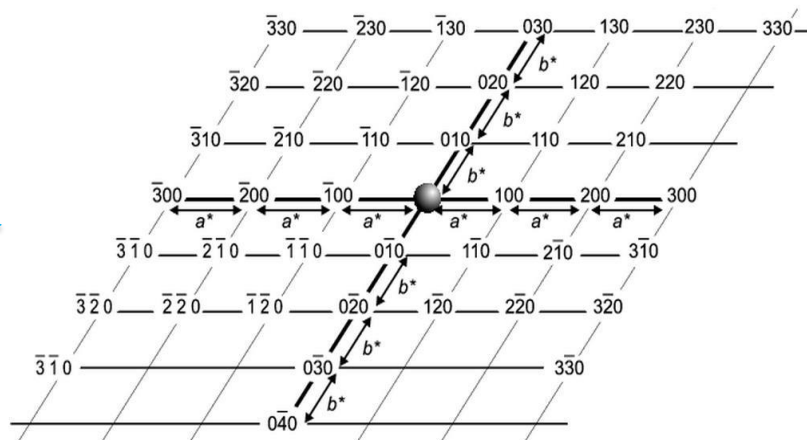
Real space



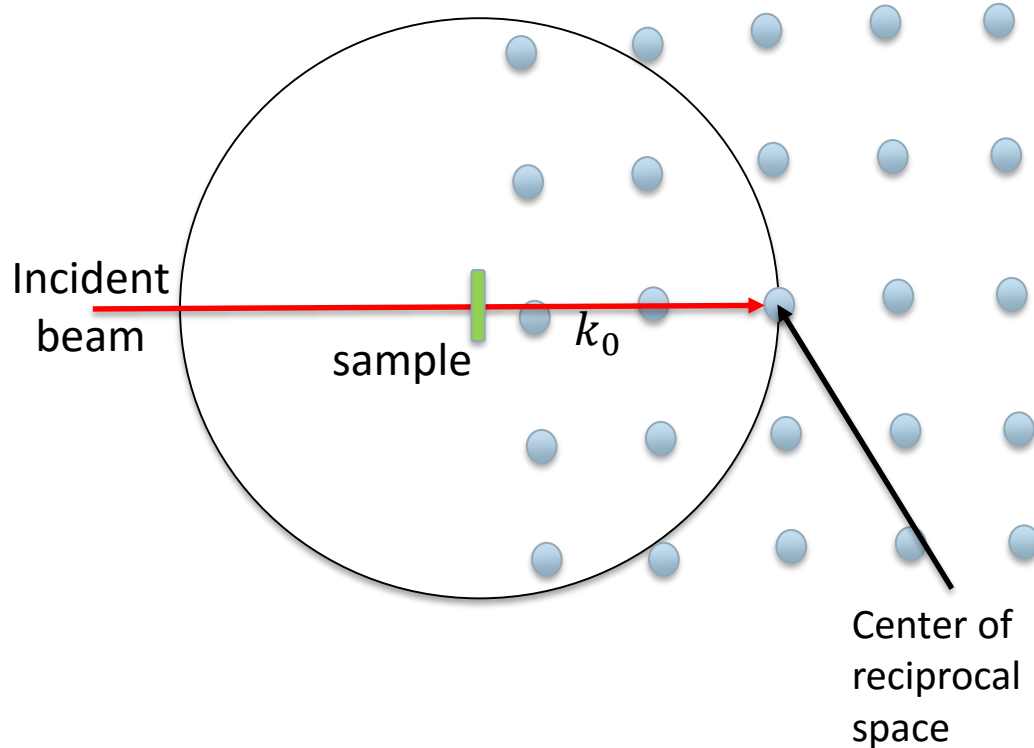
Diffraction



Reciprocal space

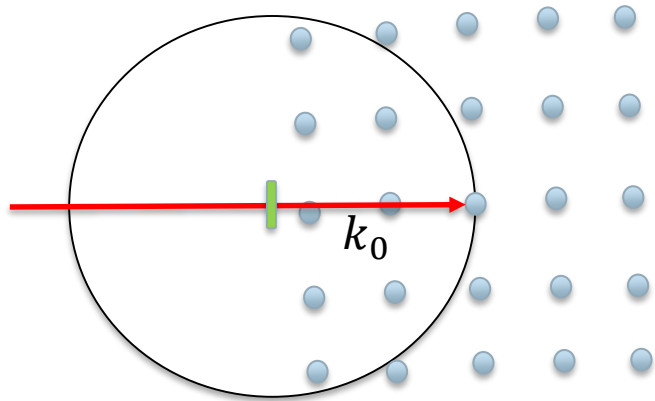


# Ewald sphere



- Ewald sphere characteristics:
- Lives in reciprocal space
  - Radius  $1/\lambda$
  - Centered in the crystal
  - Diffraction direction

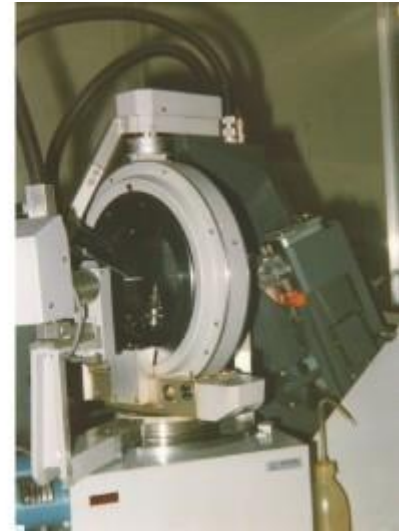
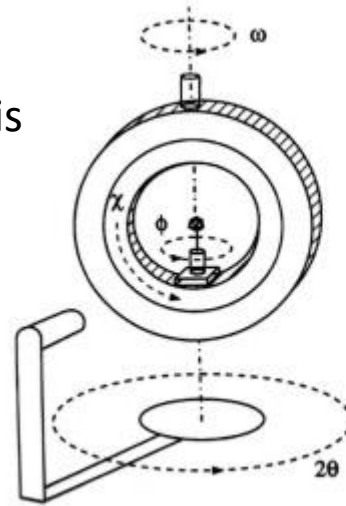
# Ewald sphere

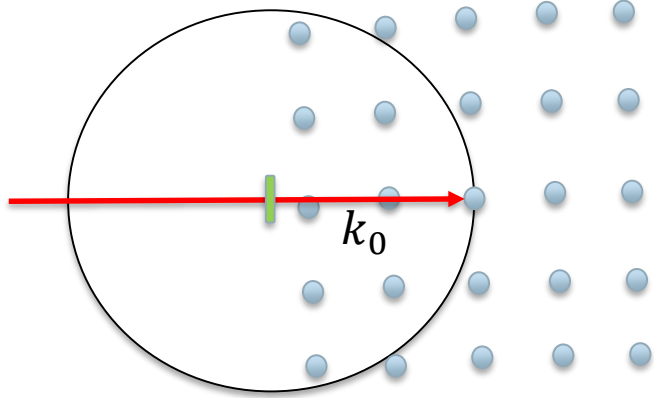


**NO REFLEXION!!**

There is not any point in the sphere's border

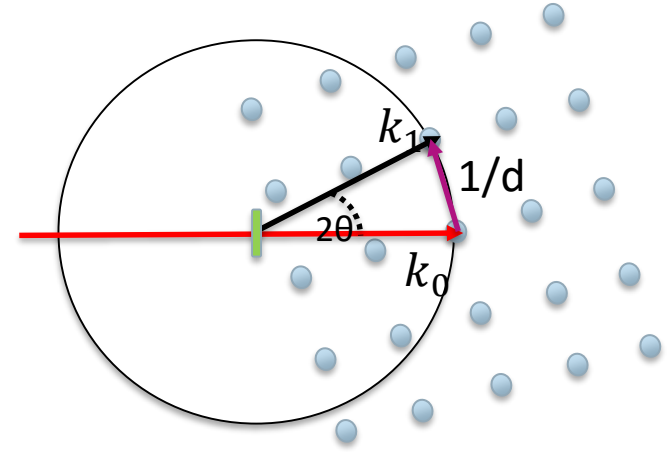
Move the real space is  
move the reciprocal  
space





### NO REFLEXION!!

There is not any point in the sphere's border

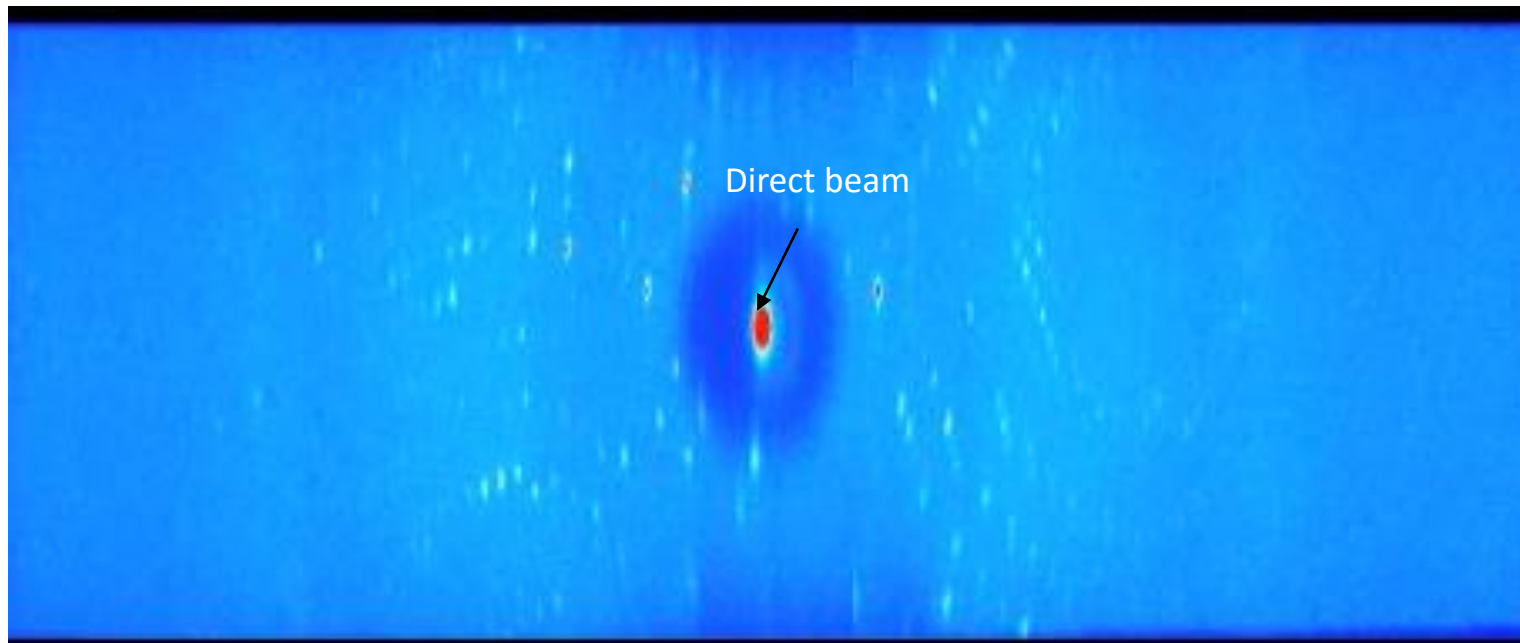


### REFLEXION!!

One point of the reciprocal space is sphere's border

$$\frac{1}{d} = d^* = ha^* = kb^* = lc^*$$

## D19 monochromatic single-crystal diffractometer

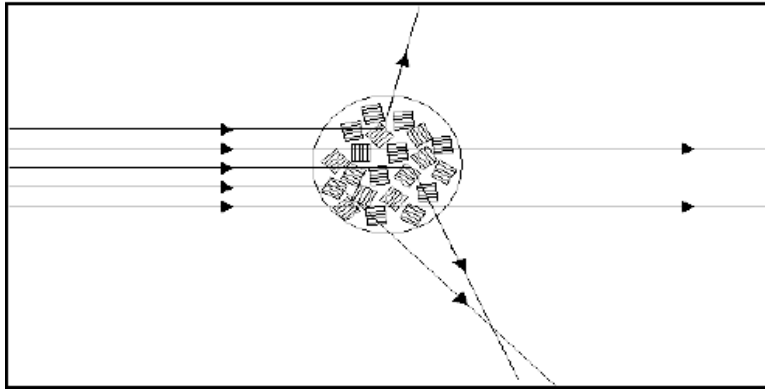




# Powder diffraction

What is powder?

Polycrystalline sample formed by a big number of single-crystals **randomly oriented**



1 cm<sup>3</sup> of powder =>  
10<sup>9</sup> (10 μm) - 10<sup>12</sup> (1 μm)  
little crystals

Always there are thousand of crystals that respect the diffraction conditions!!

# Why powder samples?

Large number of preparative methods available

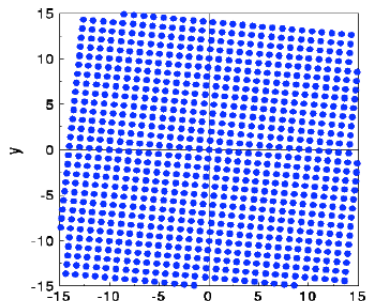
Can be prepared in large quantities (g, kg, etc.)

Fast synthesis

Real world materials are often in polycrystalline

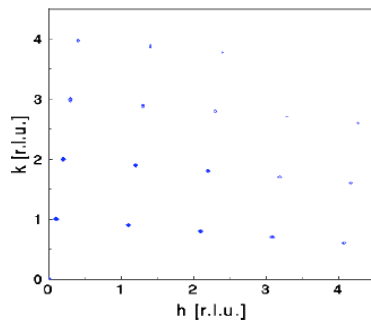


Real space

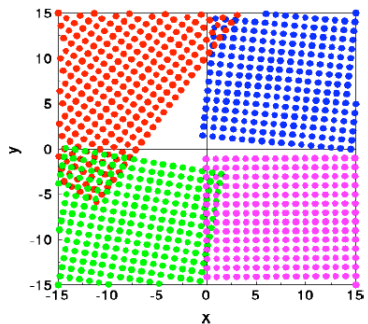
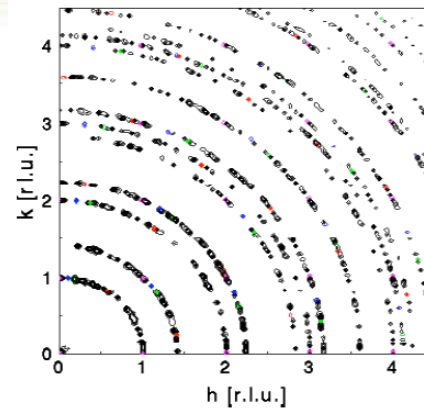


Single crystal

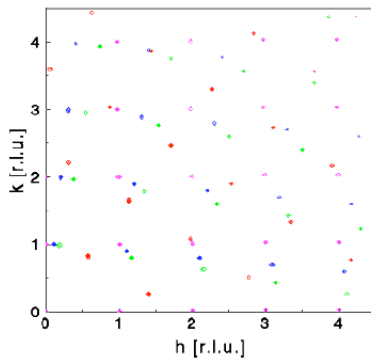
Reciprocal space



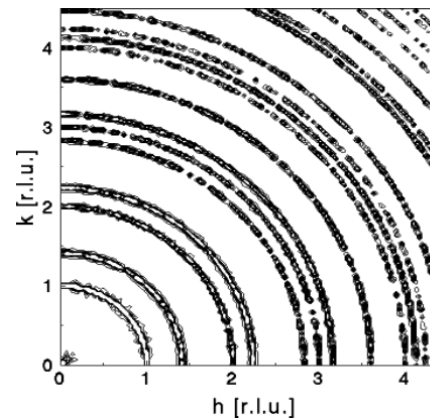
Forty single crystal



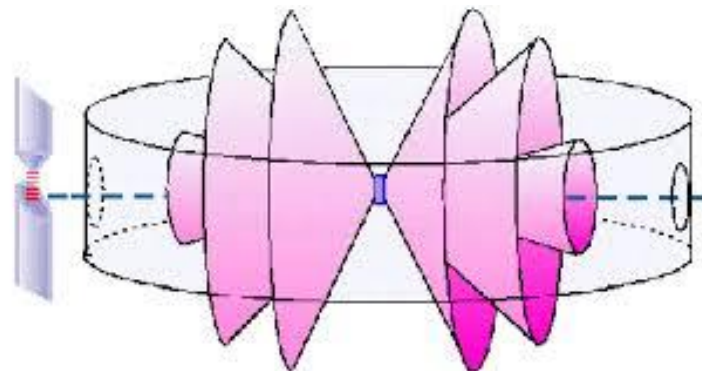
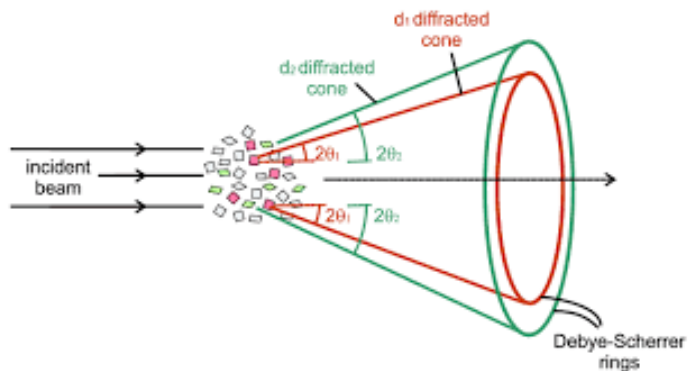
Four single crystal



Two thousand  
single crystal

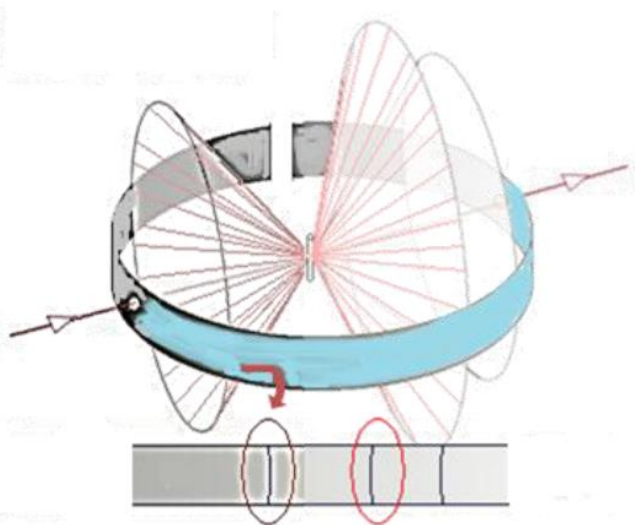


# Debye-Scherrer configuration



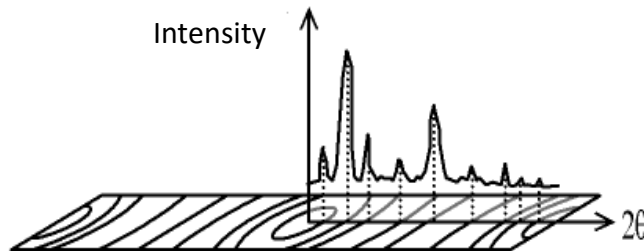
Big number of small crystal have **same orientation** => **same diffraction conditions** => reflection with **same angle**

# Debye-Scherrer configuration

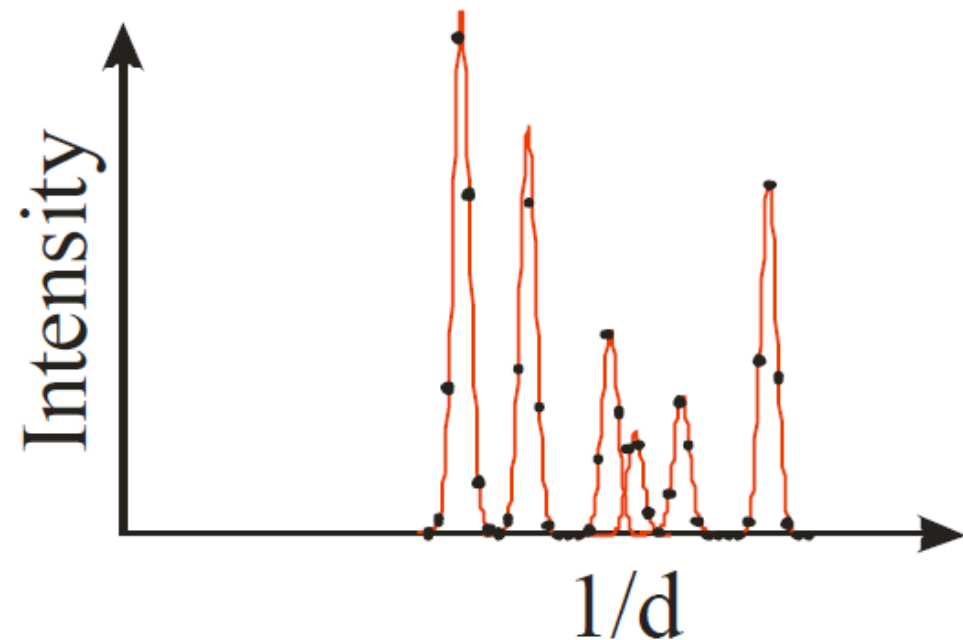


The diffracted cones intersect a banana-detector

Arcs produced in banana-detector are our diffraction pattern



## The importance of the peak shape



Rietveld method consist in use all **intensity profile data** of each reflection instead of integrated intensities

Rietveld method makes possible to obtain all the information that is buried in the superposed peaks:

- \* Space group
- \* Atom position
- \* Occupancy

# The importance of the peak shape

Two contribution to broadening:

- a) Instrumental
- b) Sample

$$I(2\theta) = \int_{-\infty}^{\infty} I_{instrument}(\psi) I_{sample}(2\theta - \psi) d\psi$$

Gaussian

$$B_{total}^2 = B_{instrument}^2 + B_{sample}^2$$

Lorentzian

$$B_{total} = B_{instrument} + B_{sample}$$

$\Psi$  is dummy variable, cover  
the full range of possibilities



## Particle size

$$\tau = \frac{K\lambda}{\beta \cos\theta} \longrightarrow \text{Scherrer equation}$$

$\tau$  is the mean size of the ordered domains, which may be smaller or equal to the grain size

$K$  is a dimensionless shape factor, with a value close to unity

$\lambda$  is the wavelength

$\beta$  is the line broadening at half the maximum intensity (FWHM), after subtracting the instrumental line broadening, in radians

$\theta$  is the Bragg angle



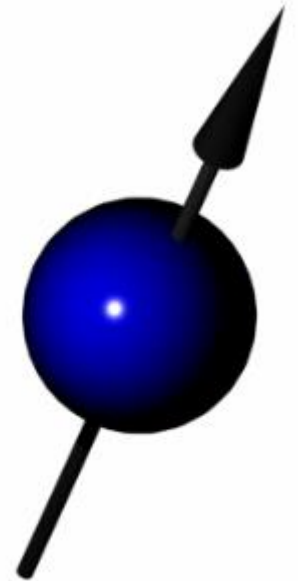
# Magnetic diffraction

## Fundamental neutron properties

- **Non-charged** particle  $m_N = 1,67 \cdot 10^{-27} kg$
- Total orbital momentum (nuclear spin)  $I=1/2$
- Neutron moment is around 960 times smaller than the electron moment

	proton	neutron
Nuclear magnetons:	$\mu_p = 2,793\mu_N$	$\mu_n = 1,913\mu_N$

For neutrons:  $\mu_n = -\gamma\mu_N\sigma$  with  $\gamma_n = 1,913$



# Magnetic diffraction

## Magnetic form factor

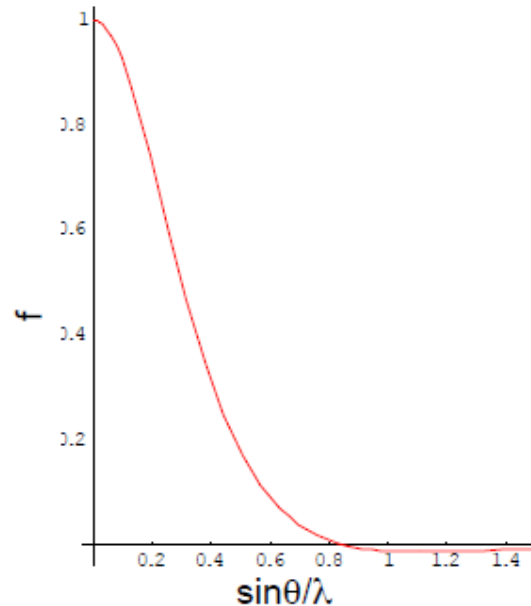
Like for nuclear scattering the Born approximation holds and the scattered amplitude is the Fourier transformation of the potential function (atomic magnetization density), the **magnetic form factor**

$$f(\mathbf{k}) = \int \rho(\mathbf{r}) \exp(i\mathbf{k}\mathbf{r}) d\mathbf{r}$$

$$f(\mathbf{k}) = \frac{g_S}{g} j_0(\mathbf{k}) + \frac{g_L}{g} [j_0(\mathbf{k}) + j_2(\mathbf{k})]$$

$g, g_S, g_L$ : g-factors

$j_n$ : spherical Bessel functions



**Magnetic information is in small Q peaks**

# Magnetic diffraction

Magnetic form factor

$$f(\mathbf{k}) = \int \rho(\mathbf{r}) \exp(i\mathbf{k}\mathbf{r}) d\mathbf{r}$$

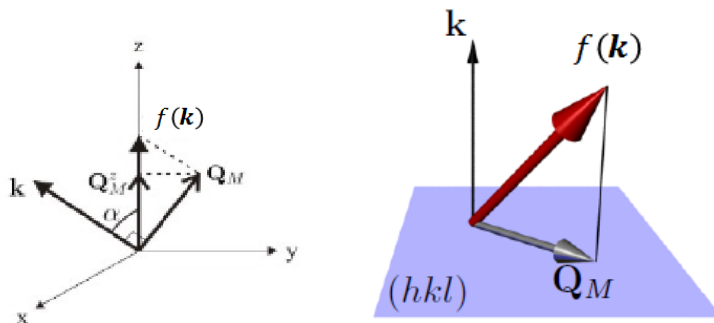
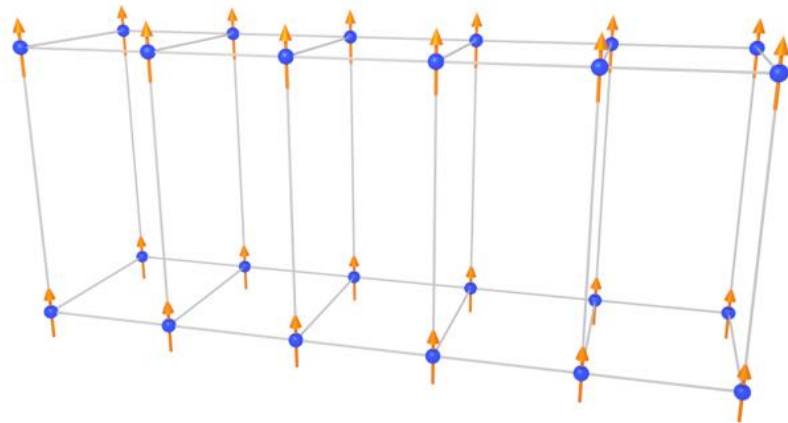
Magnetic information

- Magnetic moment
- Orientation

Information about  
the lattice

Only the component of  $f(\mathbf{k})$  which is **perpendicular** to  $\mathbf{k}$  contributes to magnetic scattering

$$Q_M = \hat{\mathbf{k}} \times (f(\mathbf{k}) \times \hat{\mathbf{k}})$$



# Magnetic diffraction

## Propagation vector

Magnetic structure **can have different periodicity and symmetry** that nuclear structure

The relation between one and another is expressed by the **propagation vector**

The propagation vector is given by **satellite peaks**

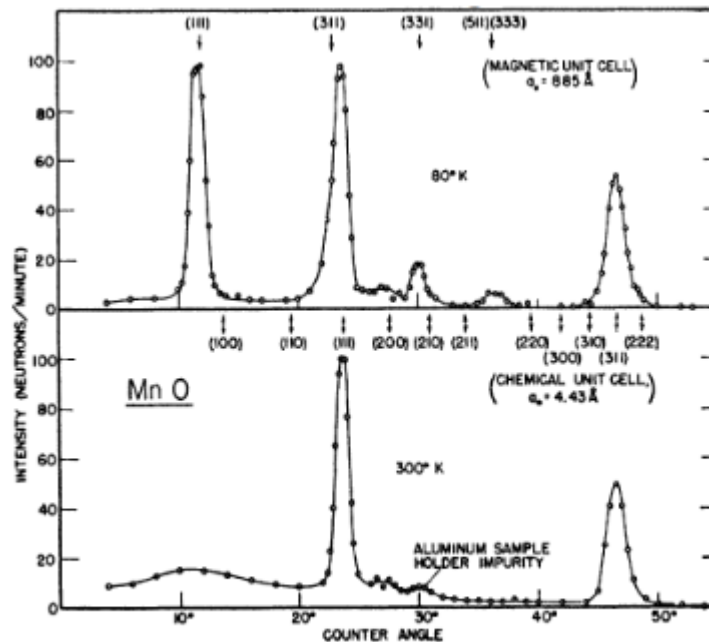
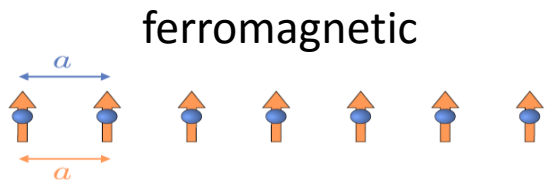


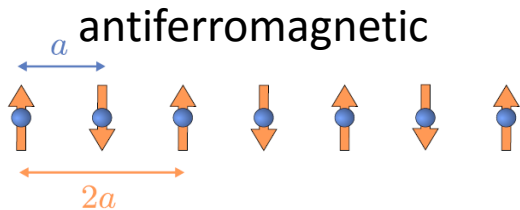
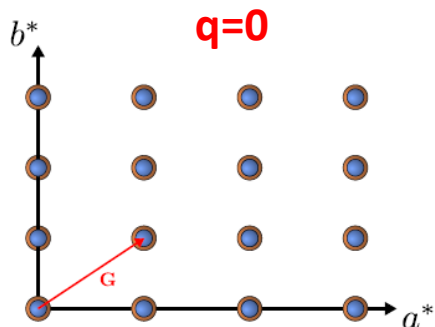
FIG. 1. Neutron diffraction patterns for MnO at room temperature and at 80°K.

# Magnetic diffraction

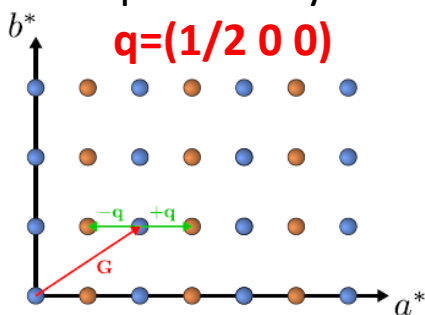
Some magnetic structures



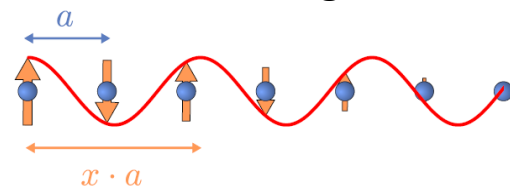
magnetic  
periodicity=nuclear  
periodicity



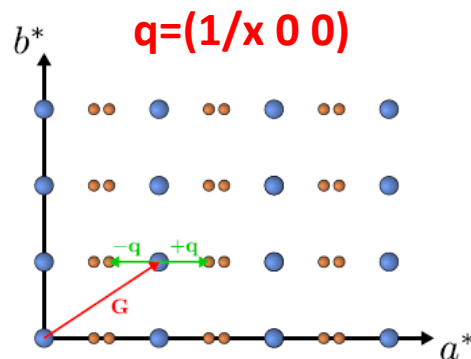
magnetic  
periodicity=2x nuclear  
periodicity



Incommensurate  
antiferromagnetic



magnetic periodicity=x  
times nuclear periodicity



# Instruments

## Single-crystal

### Nuclear structures:

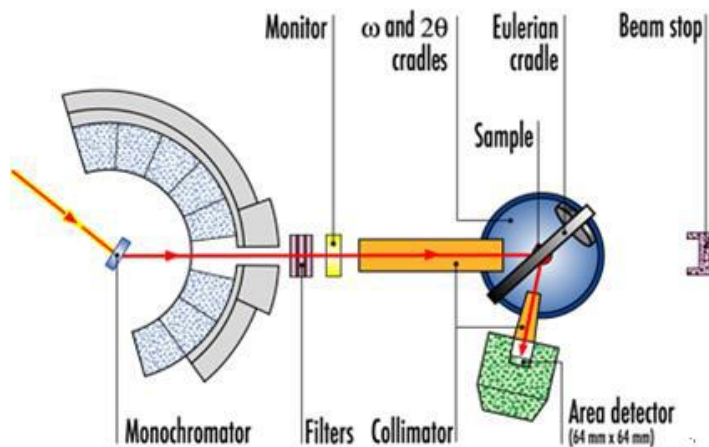
- D9
- D19
- Laue diffractometers

### Magnetic structures:

- D3
- D23
- D10

# Instruments

## D9: hot neutron 4-circle diffractometer



### Detector:

- Pixel size  $2 \times 2 \text{ mm}^2 / 0,25^\circ \times 0,25^\circ$

### Technical characteristics:

- 4-circle configuration
- Normal beam

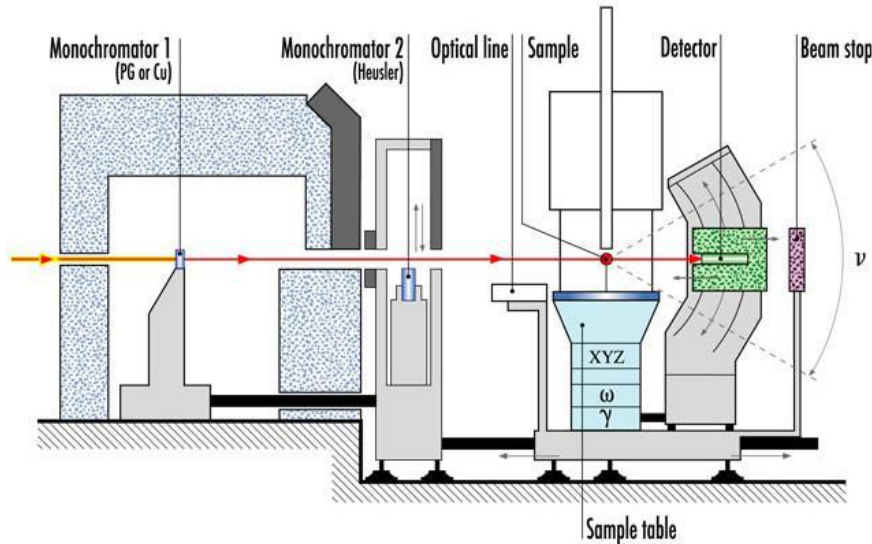
### Sample environment:

- Pressure cells
- Dilution 50mk
- Cryofurnace (1.8-520K)
- Displex and furnace (2/12K-1200K)

Inorganic compounds or  
small molecules

# Instruments

## D23: lifting-counter two-axis diffractometer



- Thermal neutron
- Either polarized or unpolarised neutron

Sample environment:

- Cryostat (1,5-300K)
- High field cryomagnet (~15T)
- Pressure cells (~30kbar)

Angular ranges:

- $-181 < \omega < 181$  °
- $-124 < \gamma < 128,5$  °
- $-28 < \alpha < 29$  °



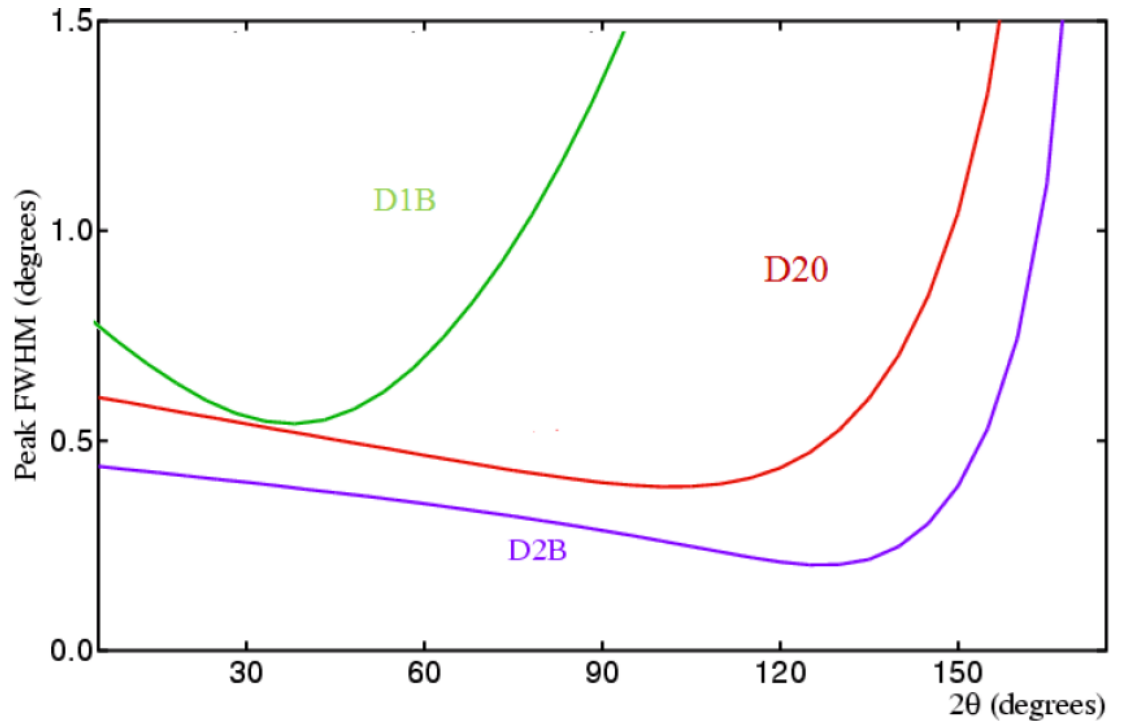
# Instruments

## Powder diffractometers

Depending of the range of better resolution, the diffractometer is useful for magnetic or for nuclear structure. The resolution depends strongly of the **take-off-angle** (diffraction angle of monochromator)

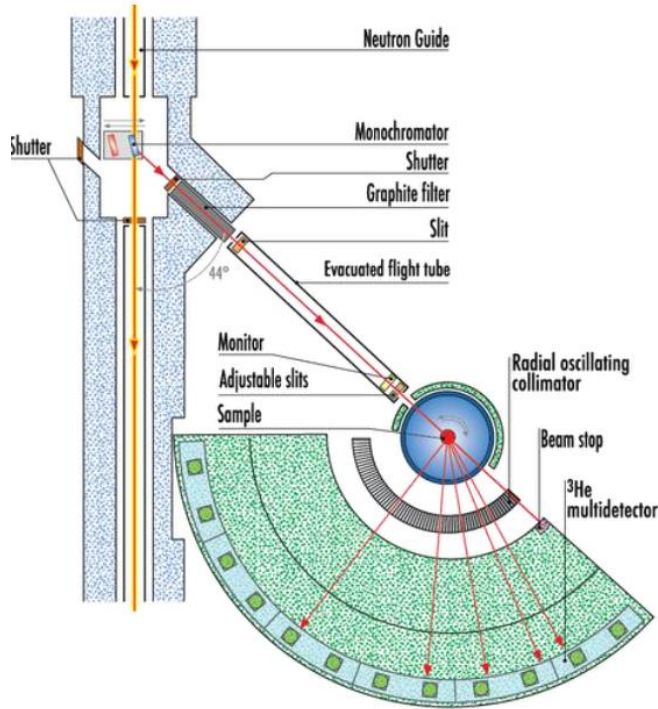
**D1B** magnetic structure

**D2B** and **D20** (high-res. config)  
nuclear structure



# Instruments

## D1B: Two-axis diffractometer



Take-off-angle:  $44.22^\circ$

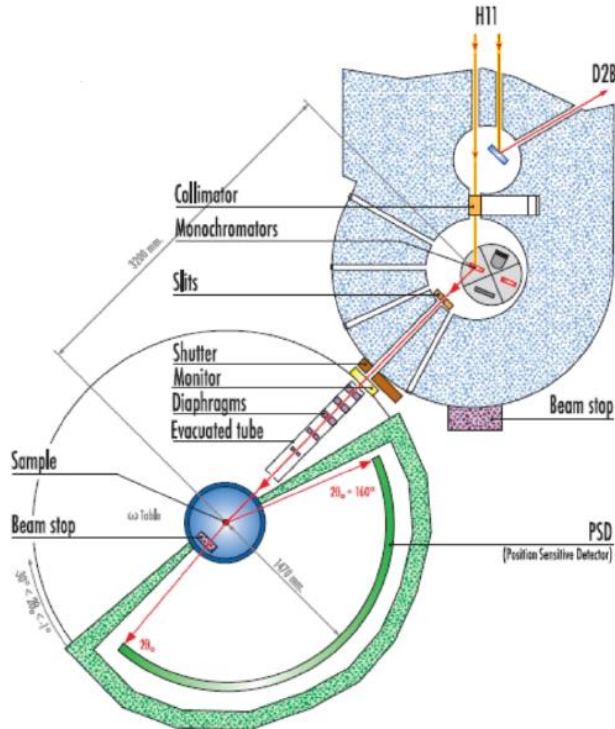
$\lambda = 1.28\text{\AA}, 2.52\text{\AA}$

Angular range sample:  $0.8^\circ < 2\theta < 128^\circ$

$^3\text{He}$  multidetector 1280 cells

# Instruments

**D20:** high-intensity two-axis diffractometer with variable resolution



5 Take-off-angle (low)= 26 °, 28 °, 30 °, 42 °

Take-off-angle (high)= 65 °, 90 °, 120 °

$\lambda=0.82-0.94\text{\AA}$ ,  $1.30\text{\AA}$  and  $2.41\text{\AA}$

$^3\text{He}$  microstrip gas-detector (PSD) 1536 cells

# Neutron or X-rays?

	X-Rays	Neutron
Identify phase	YES	YES
Space group	YES	YES
Structure refinement	YES	YES
Light elements	NO	YES
Neighbouring elements	NO	YES
Structure solution	YES	YES
Magnetic structures	NO	YES



*INSTITUT LAUE LANGEVIN*

# Magnetic cross section

For elastic neutron magnetic scattering, is necessary evaluate the cross section

$$\frac{d\sigma}{d\Omega} = \left( \frac{m}{2\pi\hbar} \right)^2 \sum p_i |\langle \mathbf{k}_f, \chi_f | V_m | \mathbf{k}_i, \chi_i \rangle|^2$$

- $\chi_i$  initial spin-state of the neutron
- $\chi_f$  final spin-state of the neutron
- $\mathbf{k}_i$  incident wave-vector
- $\mathbf{k}_f$  scattered wave-vector
- $\mathbf{Q}$  momentum transfer
- $V_m$  magnetic interaction potential
- $p_i$  : probability to find the initial spin-state  $i$

