

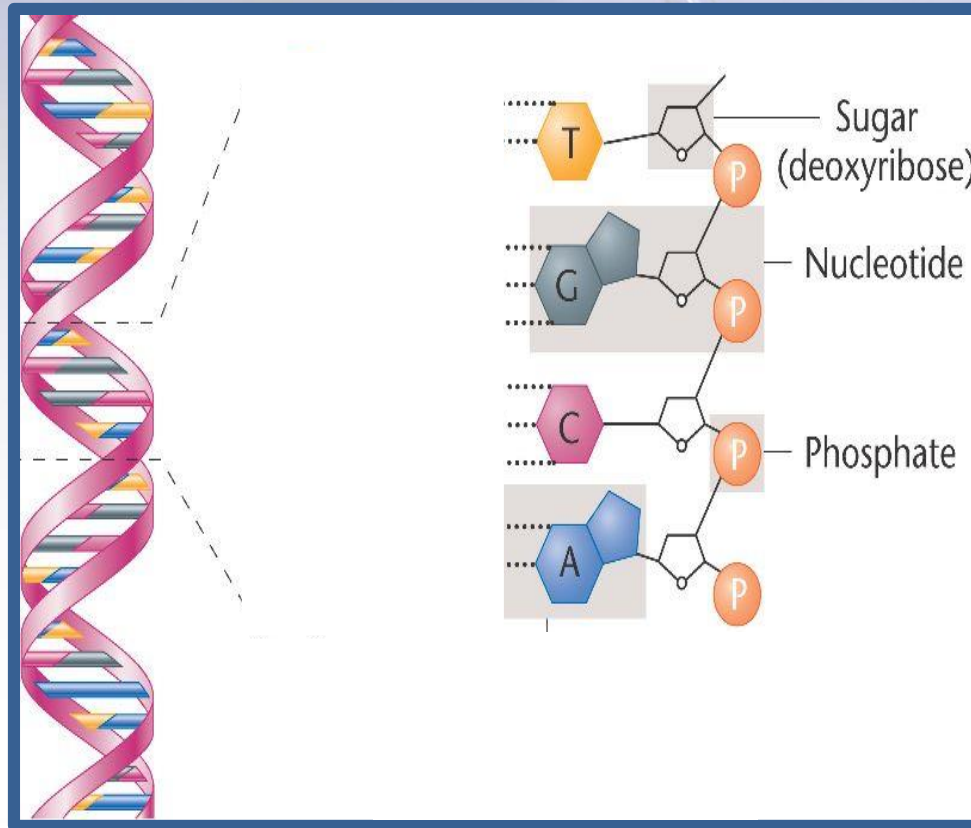


Spatial correlations in the melting transition of DNA

All you need is neutrons ILL seminars 27 October 2015

Introduction

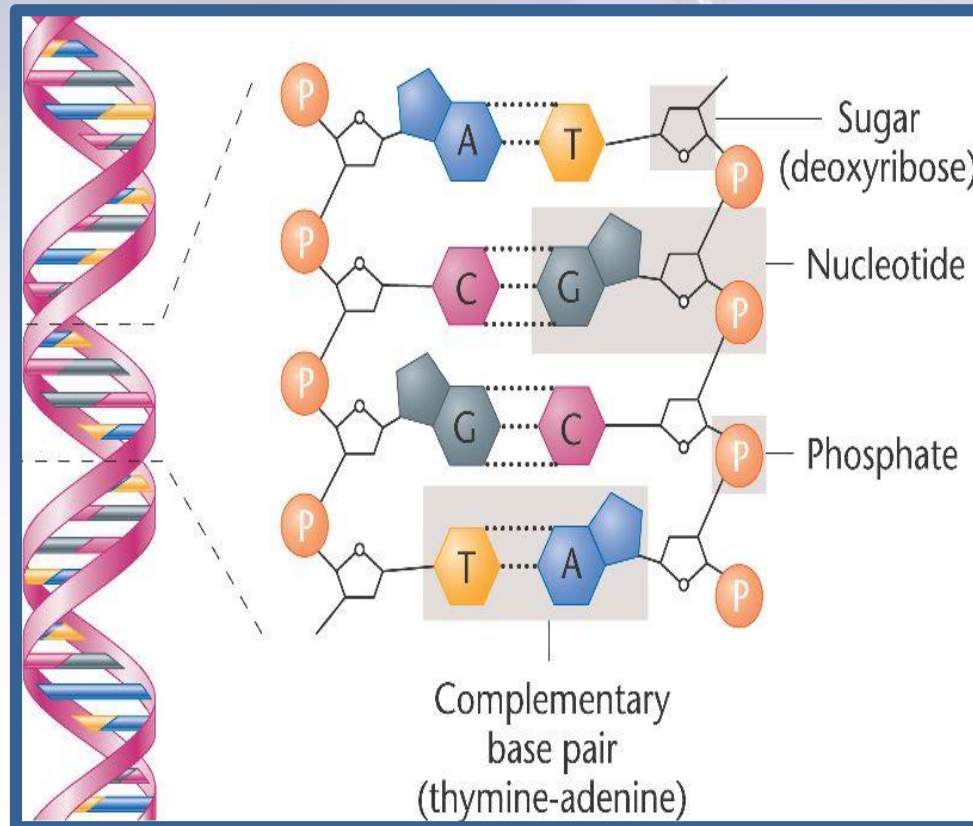
DNA STRUCTURE:



www.cliparthut.com

Introduction

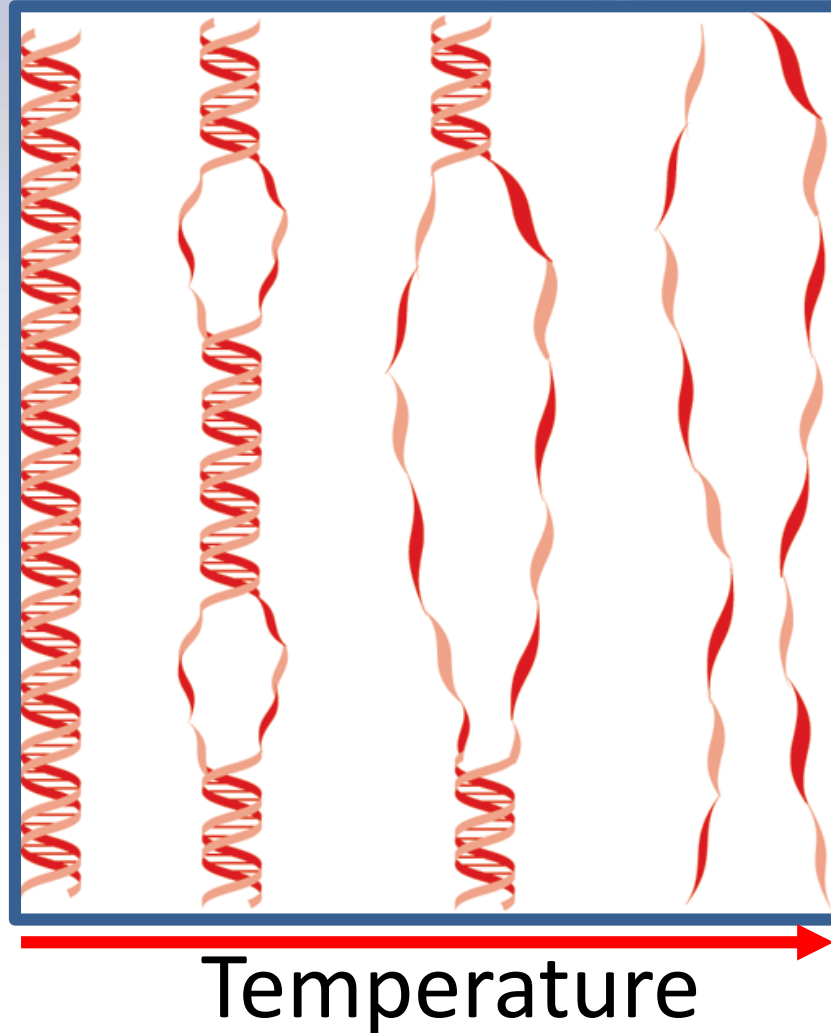
DNA STRUCTURE:



www.cliparthut.com

Introduction

DNA MELTING:



M. Peyrard, *Nature Physics* 2006, 2, 13–14.

Introduction

Why melting transition?

Applications:

- Polymerase Chain reaction.
- High resolution melting analysis.

Theory:

- 1D phase transition.

What's original in our approach?

Previous techniques:

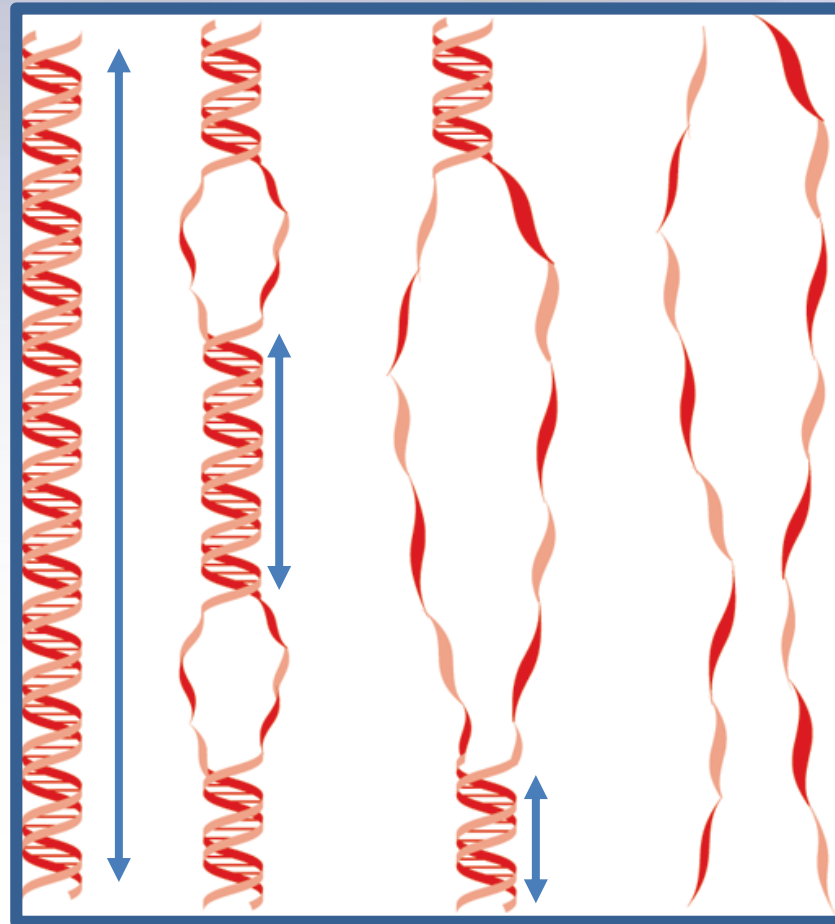
- DSC.
- Fluorescence.
- NMR.
- UV-Vis absorption.

Neutrons:

- Spatial correlation.

Introduction

CORRELATION LENGTH:



M. Peyrard, *Nature Physics* 2006, 2, 13–14.

Temperature

- ❖ The correlation length is related with the average size of the closed domains during the melting.
- ❖ Size of the open domains can be infer.

Methods

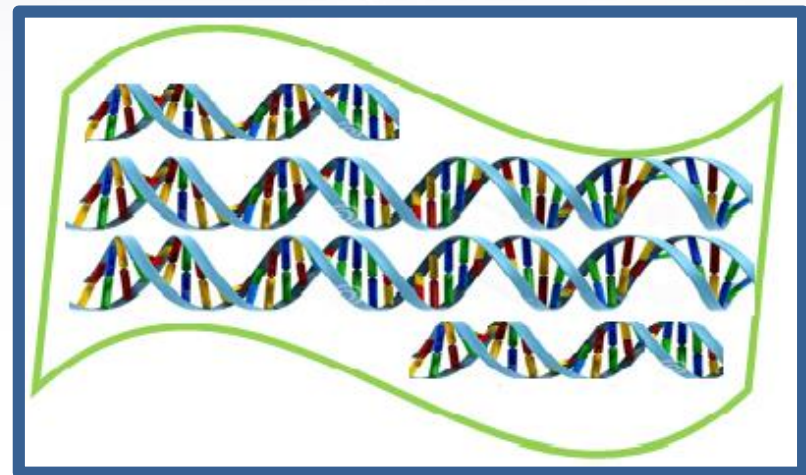
- ❖ Calorimetry, X-ray diffraction.
- ❖ Neutron diffraction:

Fiber DNA \longrightarrow Semi-crystal \longrightarrow Bragg Peaks

DNA film samples:

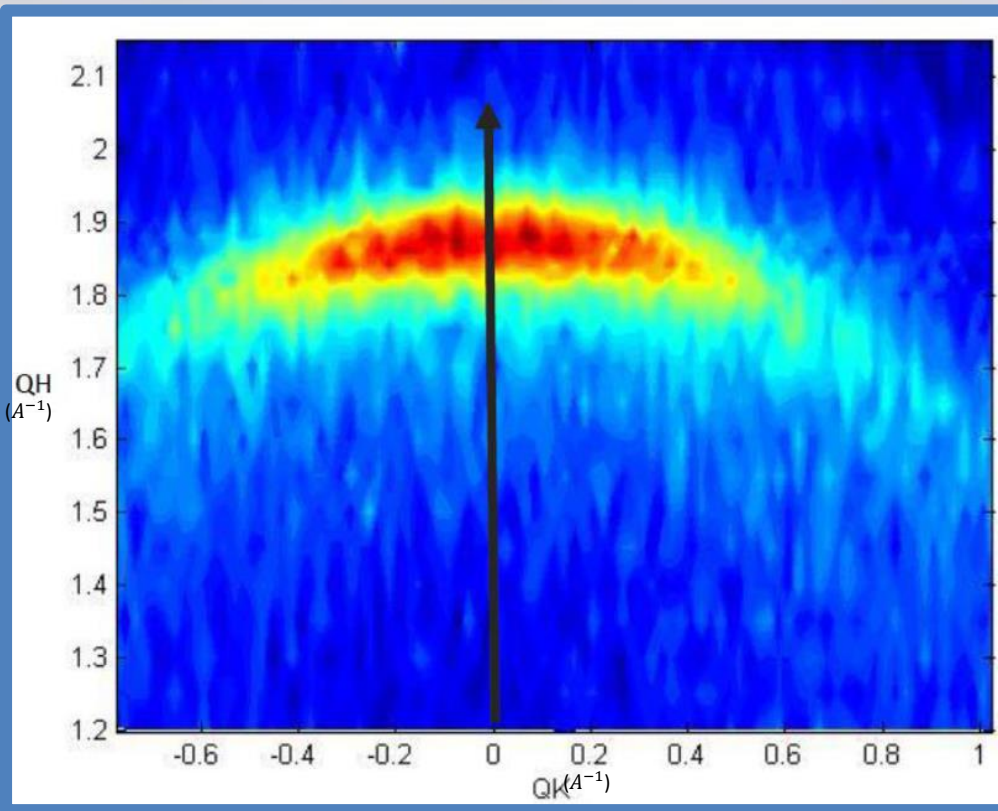


DNA-orientation in the film:



Methods

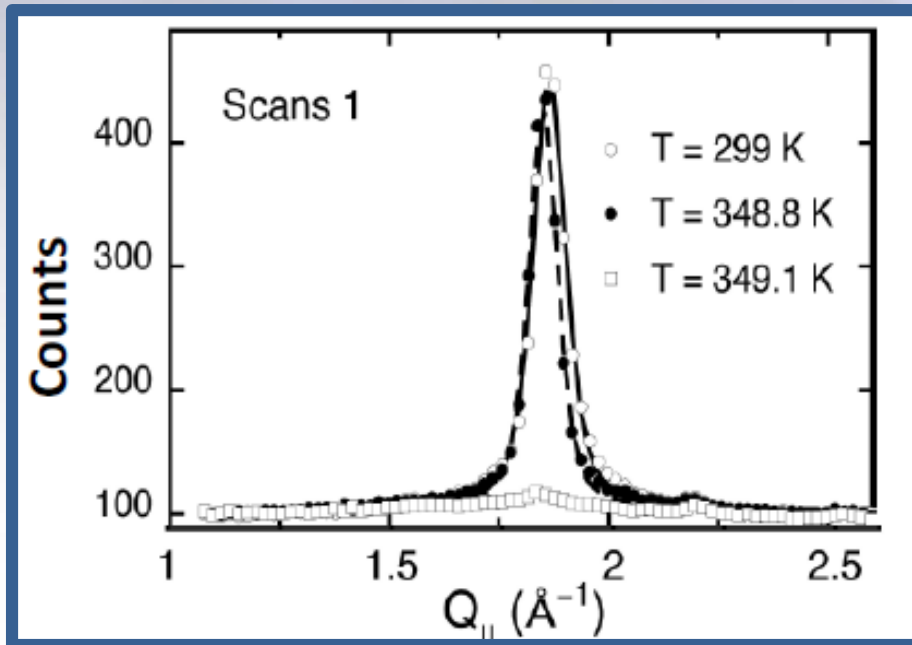
Fiber DNA \longrightarrow Semi-crystal \longrightarrow Bragg Peaks



- ❖ $Q_H = 1.85 \text{ \AA}^{-1}$
 $\rightarrow \text{spacing} = 3.4 \text{ \AA}$
- ❖ Related with the correlation of two closed base-pairs.
- ❖ The scattering elements which are the origin of this peak are **closed base pairs**.

Dry Fiber

Longitudinal cut of the Bragg peak:



Wildes, A.; Phys. Rev. Lett. 2011, 106, 048101.

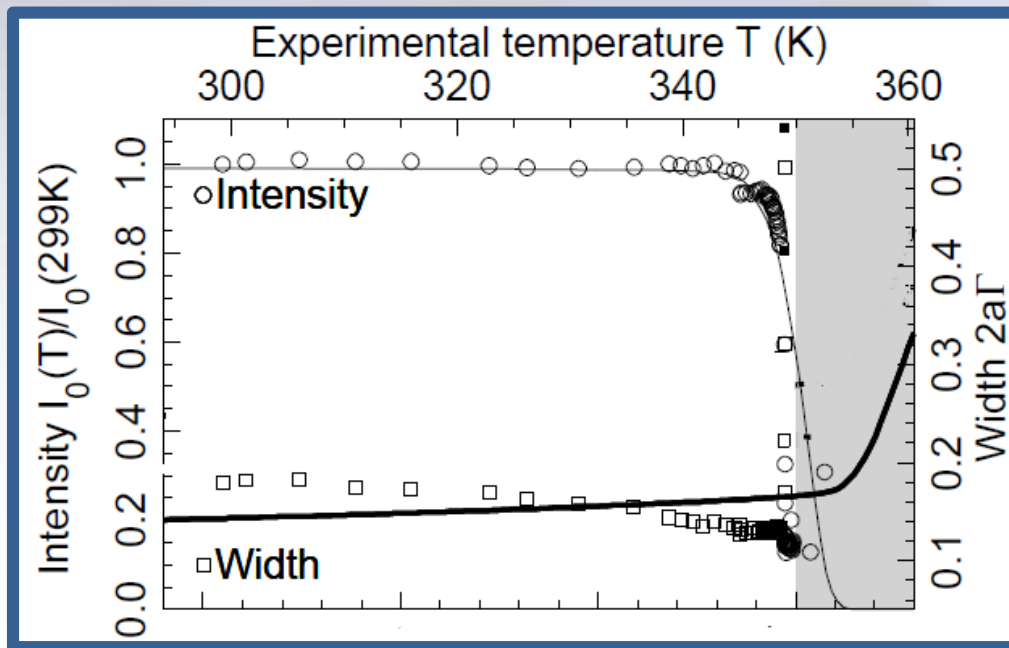
❖ I_0 (integrated I) $\propto n^0$
of closed base pairs.

❖ $\Gamma \propto \frac{1}{\text{correlation length}}$

❖ If $T \uparrow$ thus $I \downarrow$ and $\Gamma \uparrow$

The Γ give us qualitative info about the evolution of the bubbles!

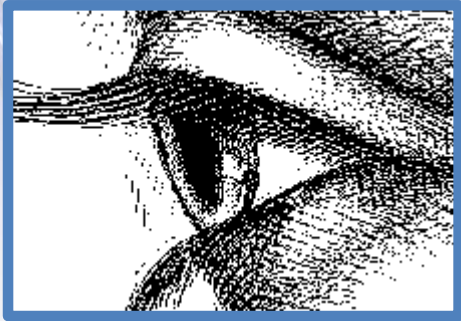
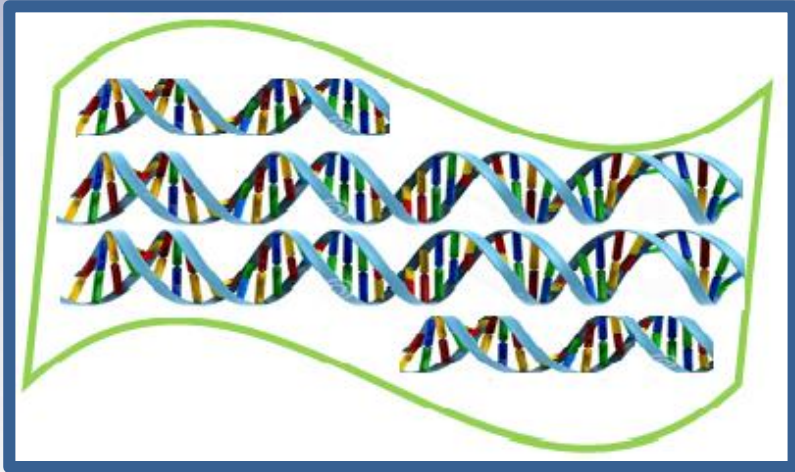
Dry Fiber



Wildes, A.; Phys. Rev. Lett. 2011, 106, 048101.

- ❖ Theory allows us to calculate the distribution of the open base pairs.
- ❖ Discrepancy of less than 15%.
- ❖ Problem: The model studies **independent molecules of DNA.**

Osmotic pressure

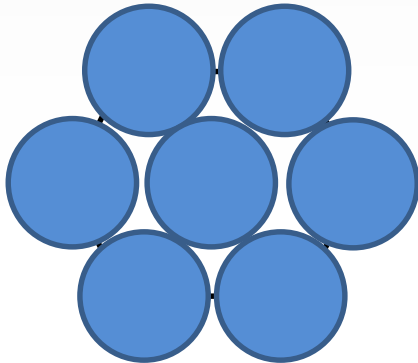


<http://www.ralphmag.org>

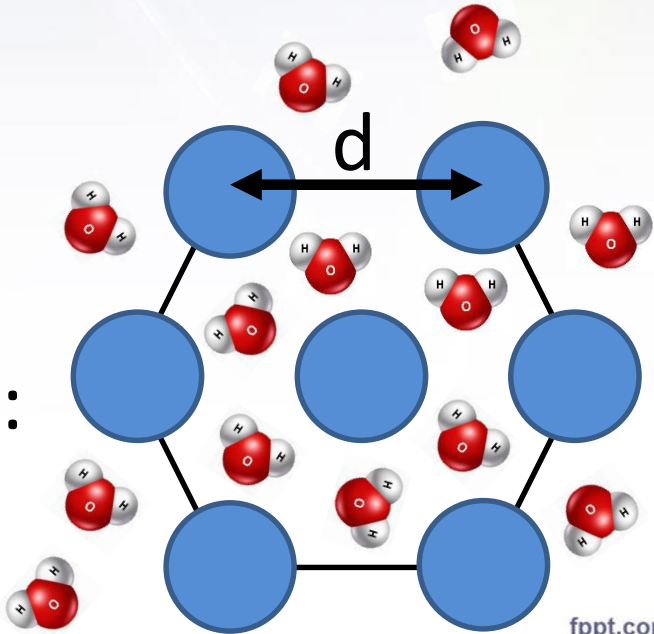


● = DNA molecule

DRY film:

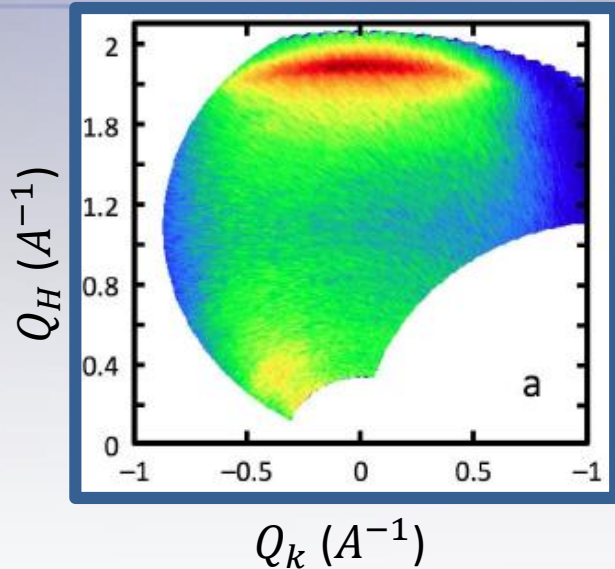


Wet fiber:

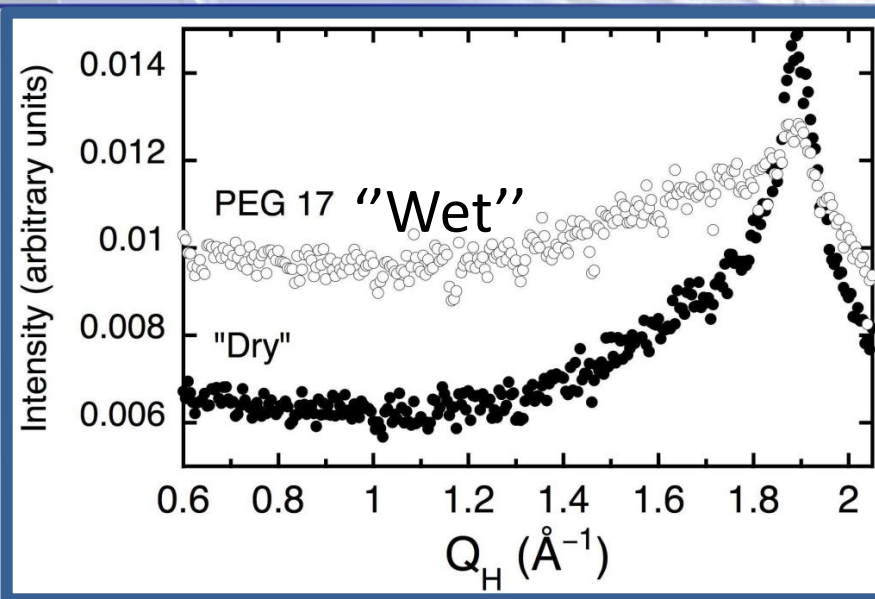
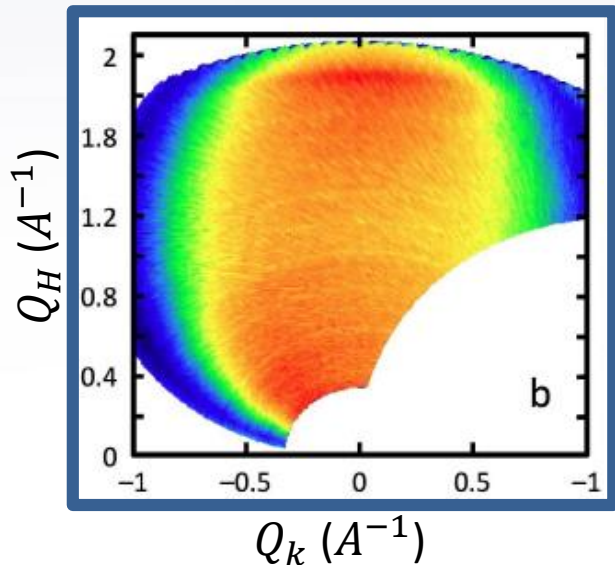


Wet Fiber

DRY



WET



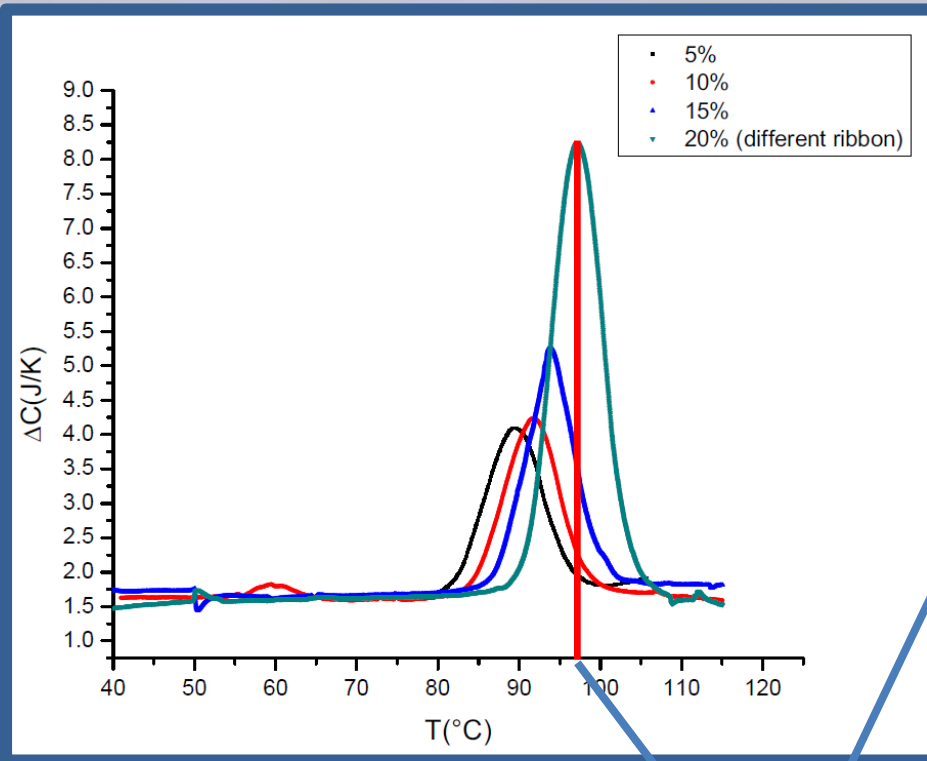
J. Phys. Chem. B, 2015, 119 (12), pp 4441–4449

Problems

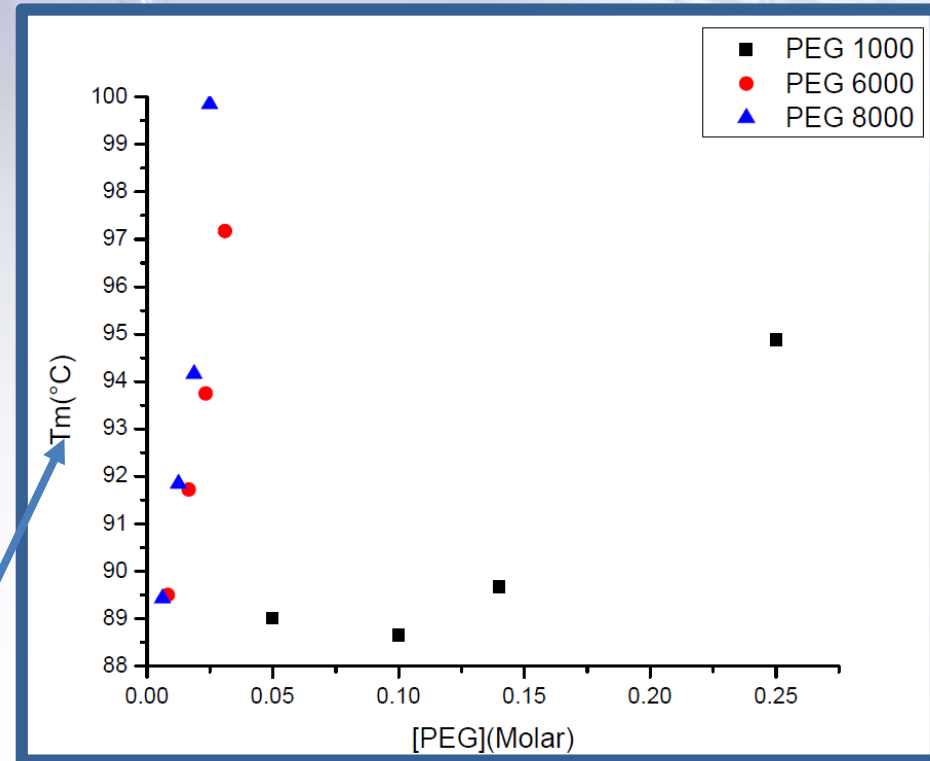
- ❖ High background.
- ❖ Peak anisotropy.
- ❖ Impossibility of a complete melting.

Calorimetry

Melting curves:



T_m vs. PEG concentration:

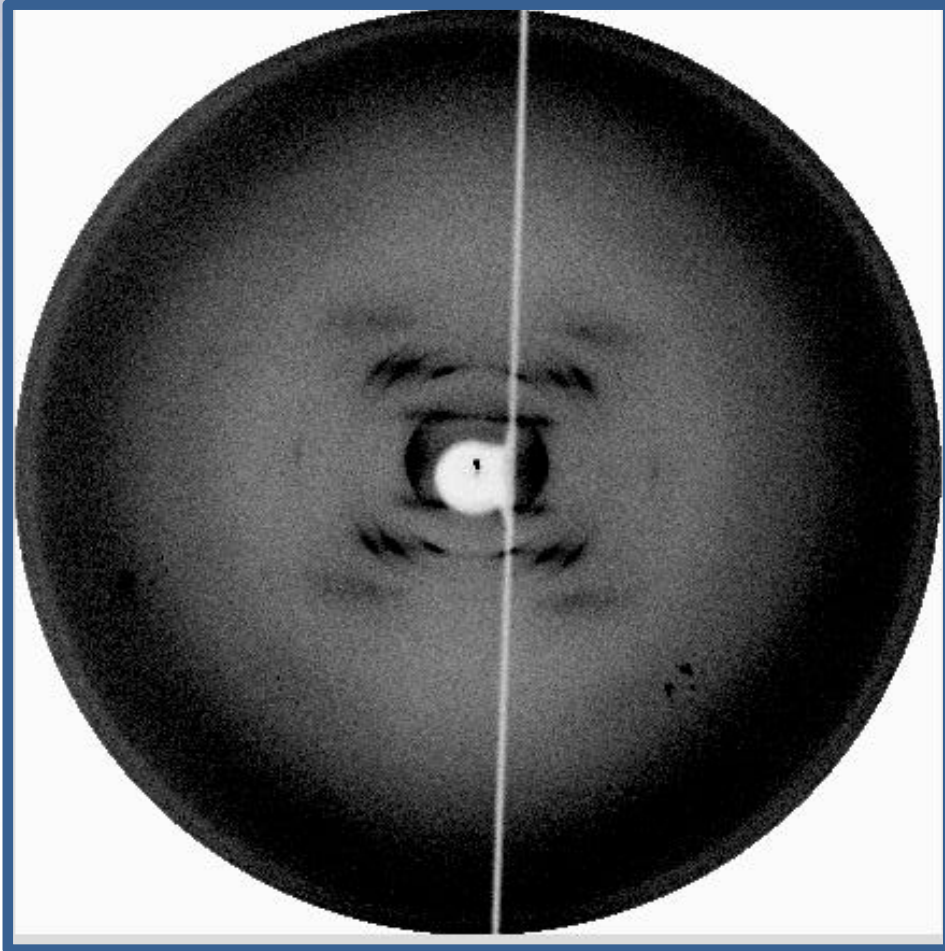


Melting temperature (T_m): Middle point of the transition.

Using a shorter PEG is better!

X-Ray diffraction

DNA in 66% Ethanol solution:

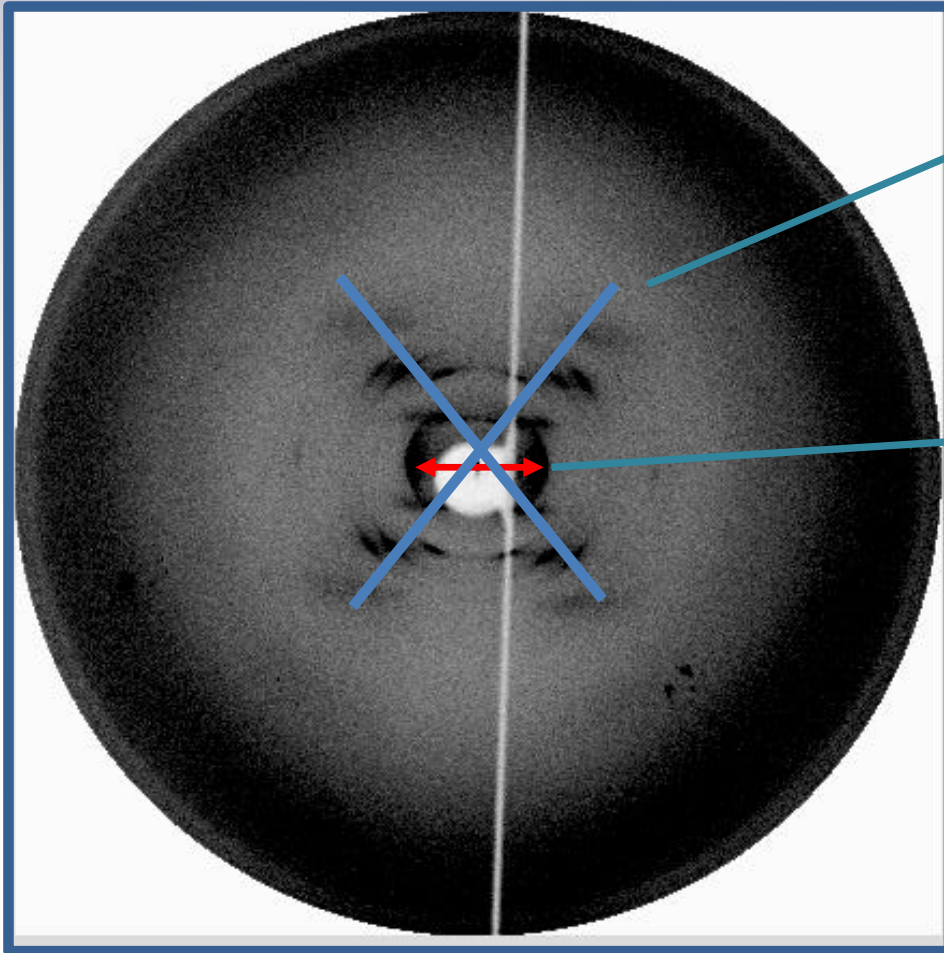


When we put the DNA in ethanol/water mixtures....

- ❖ ...will it still be crystalline?
- ❖ ...will the intermolecular space grow?

X-Ray diffraction

DNA in 66% Ethanol solution:



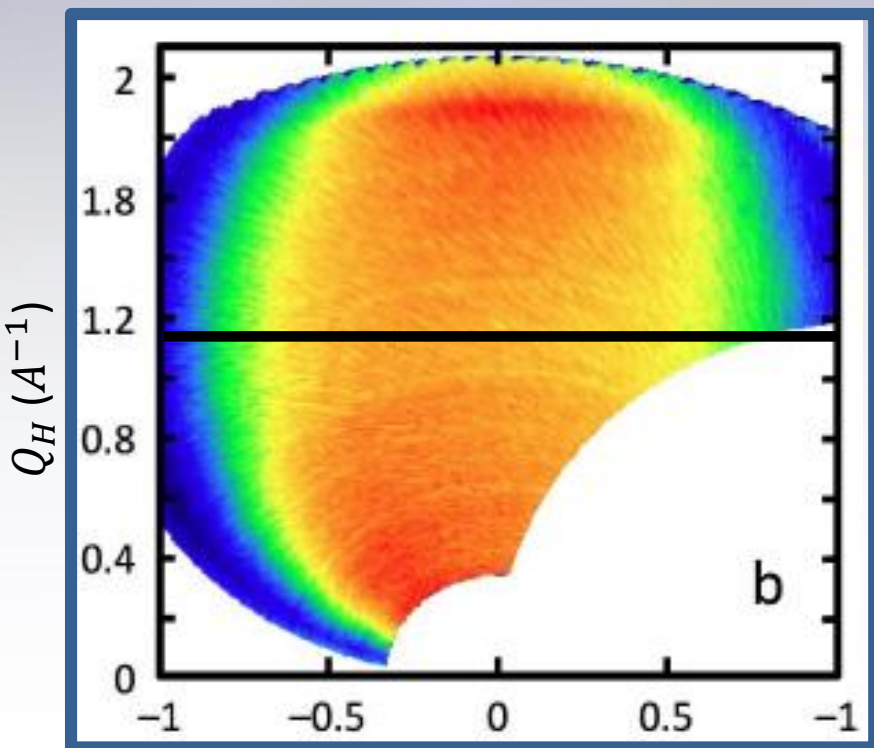
The X pattern is a sign of the semicrystalline B form.

Equatorial spacing $\propto \frac{1}{d}$

d =intermolecular spacing

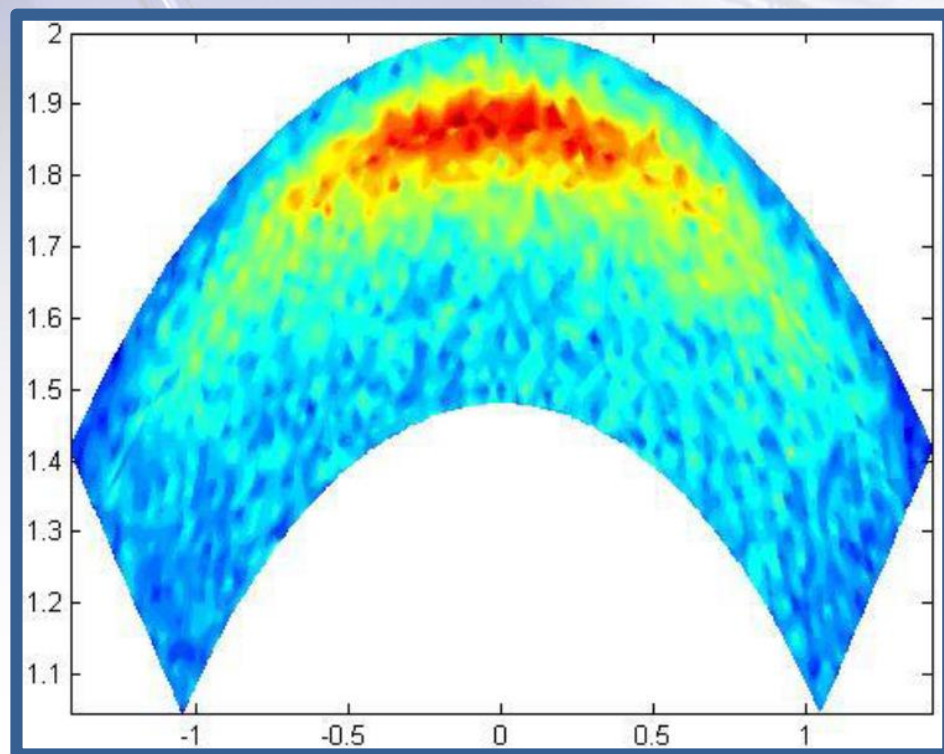
Wet Fiber 2

Before



J. Phys. Chem. B, 2015, 119 (12), pp 4441–4449

After

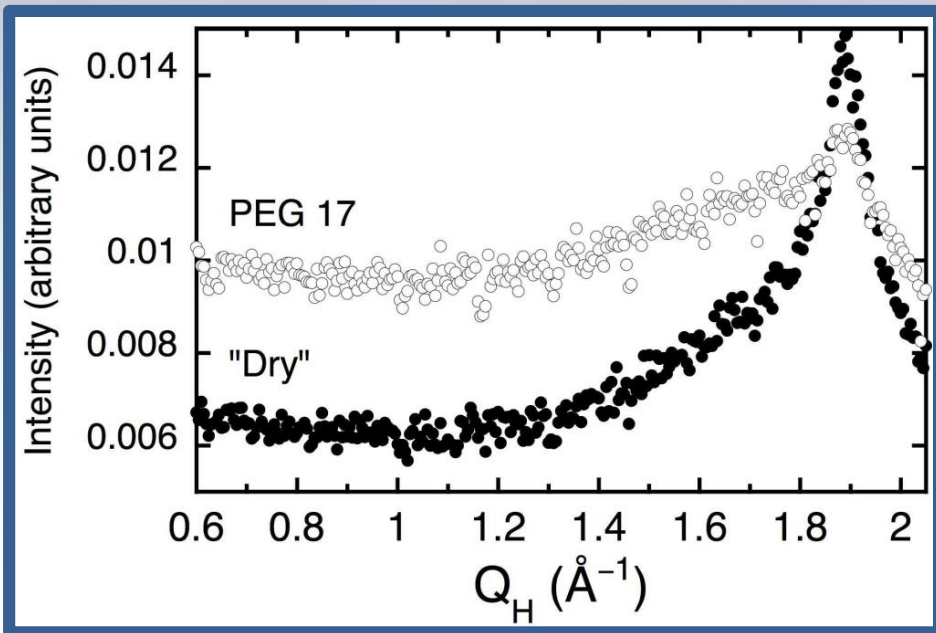


Q_k (\AA^{-1})

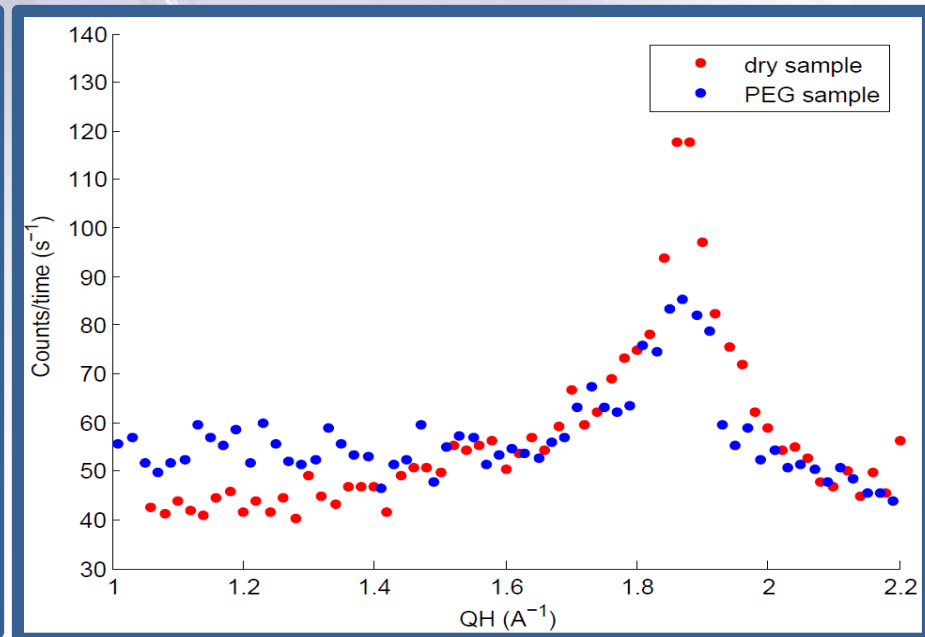
Better ratio peak-signal to noise!

Wet Fiber 2

Before



After



More symmetric peak!

Conclusions

- ❖ The melting transition in DNA is not fully understood due to the lack of spatial information.
- ❖ Neutrons can account for this lack.
- ❖ Osmotic pressure methods allow us to study the melting of fiber DNA in solution and prove the effect of the confinement.

Acknowledgment

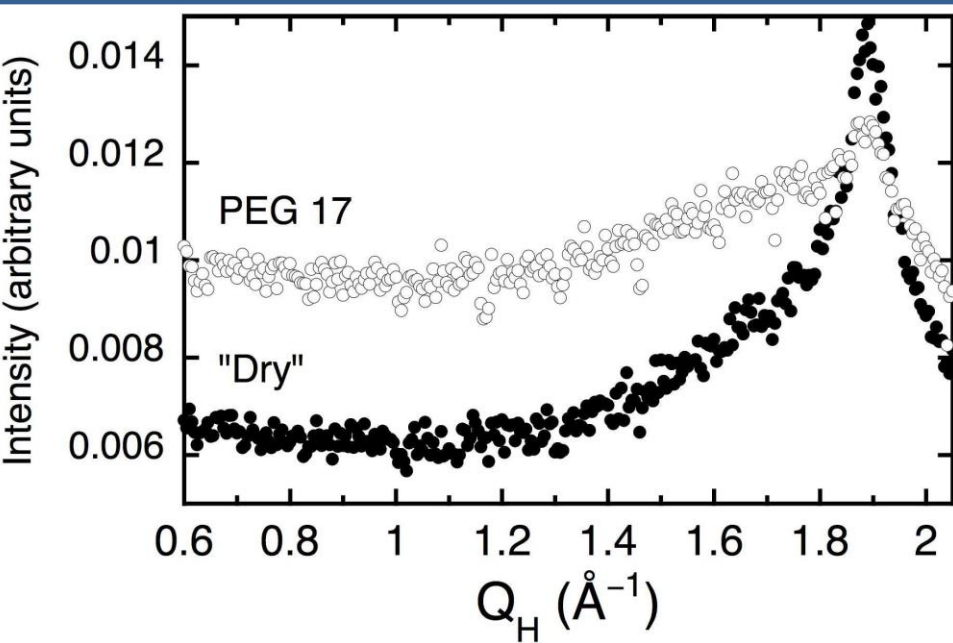
Thank you for your attention

Andrew Wildes, Santiago Cuesta López, Michel Peyrard, Jean-Luc Garden, Gael Moiroux, David Hess, Matthew Blakeley, Diane Lançon, Hassan Belrhari, Eva María Villar Alvarez, Andrew MC Carthy...

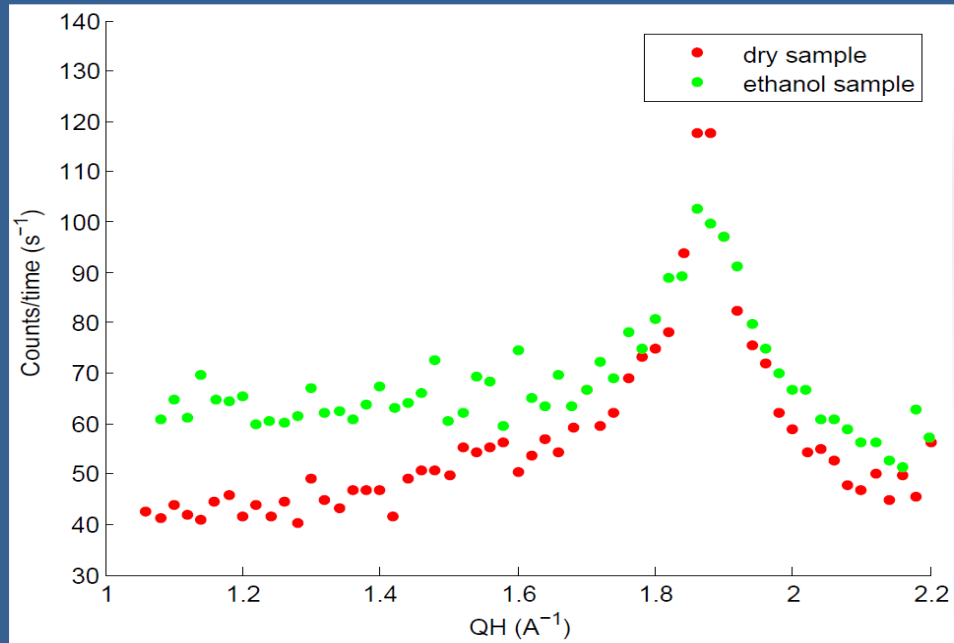


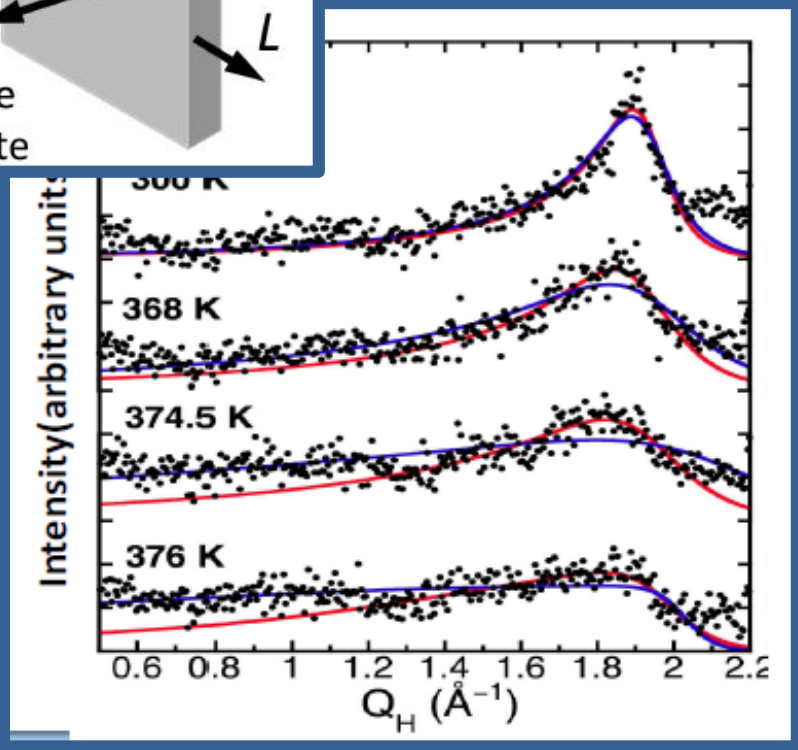
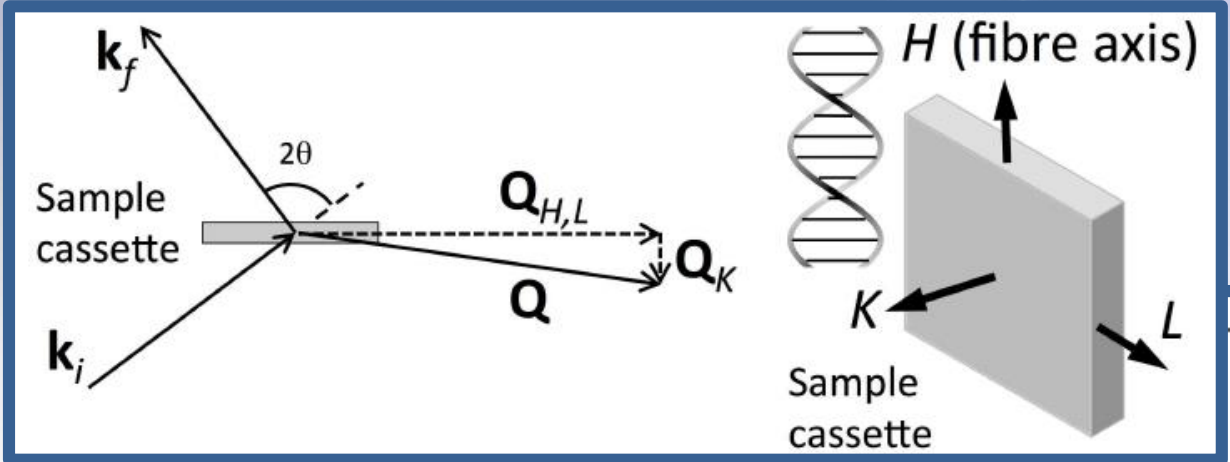
Wet Fiber

Before

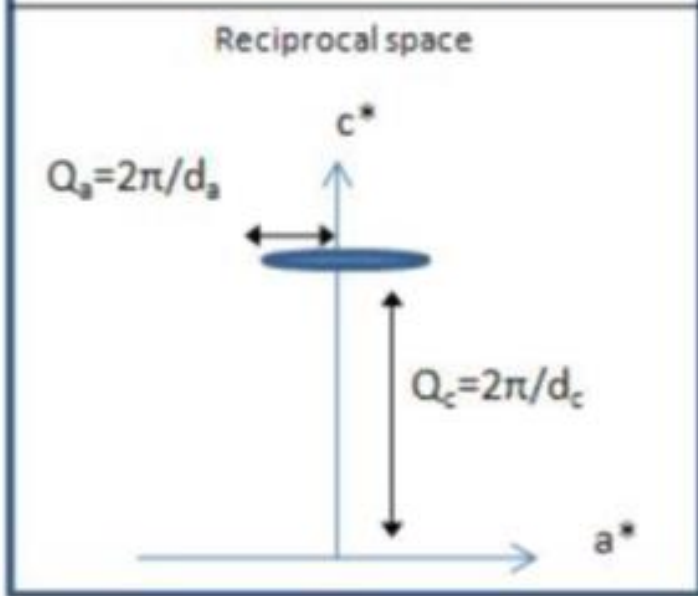
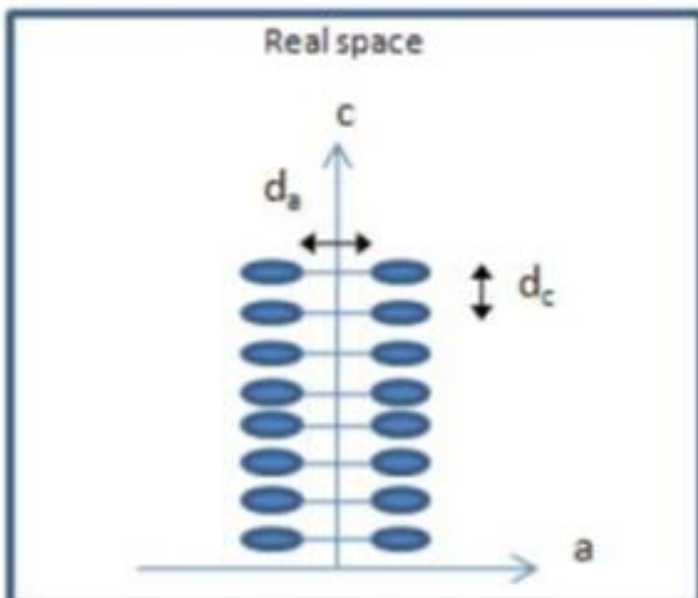


After

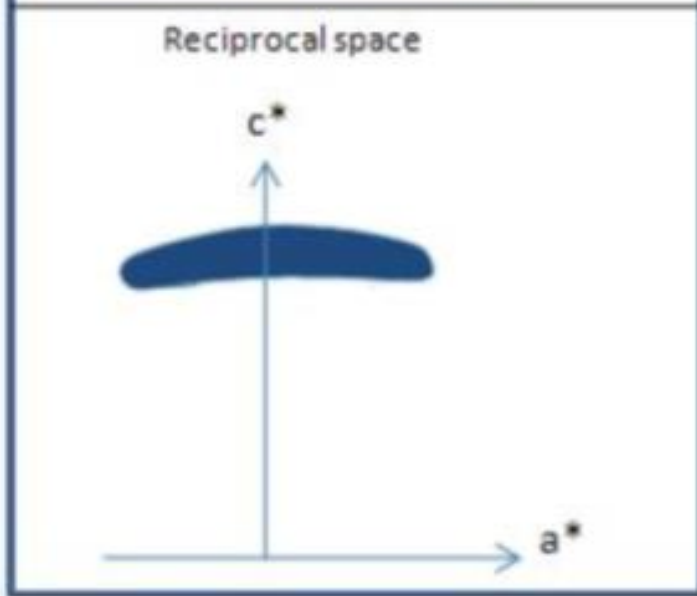
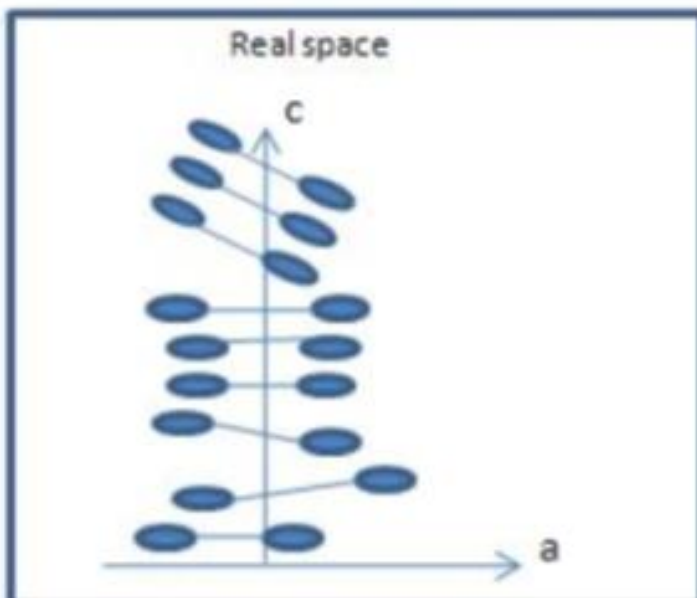


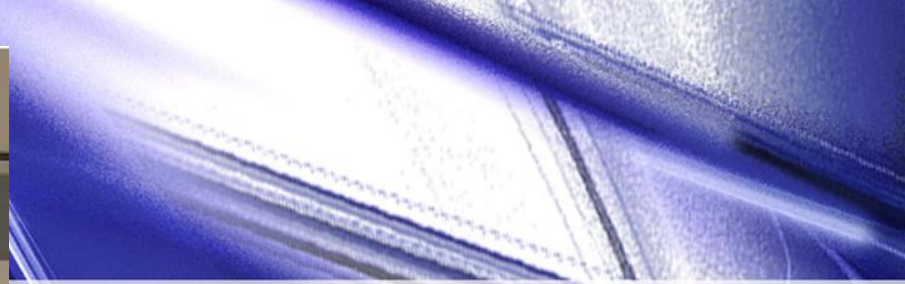


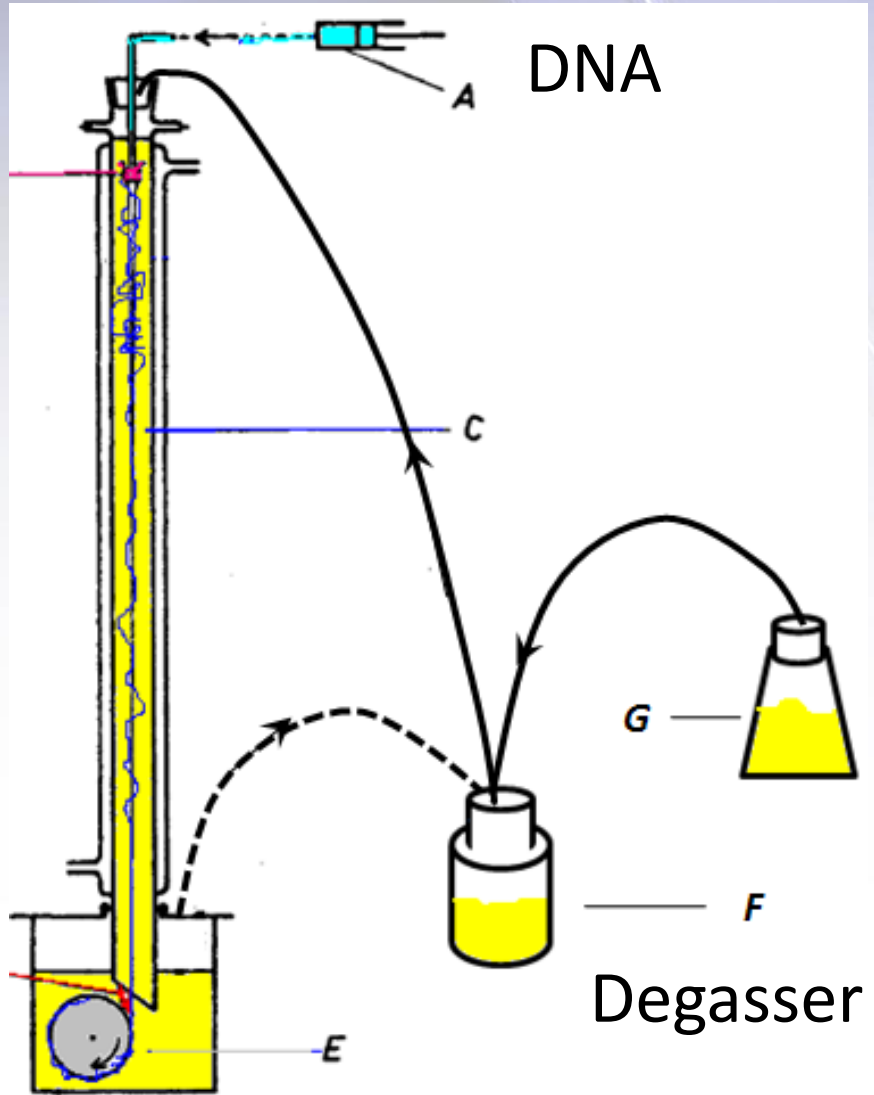
Perfect 1-D crystal



Fluctuations in c- and a- direction and mosaicity







DNA

Ethanol

Degasser

Neutrons

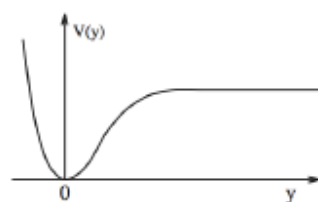
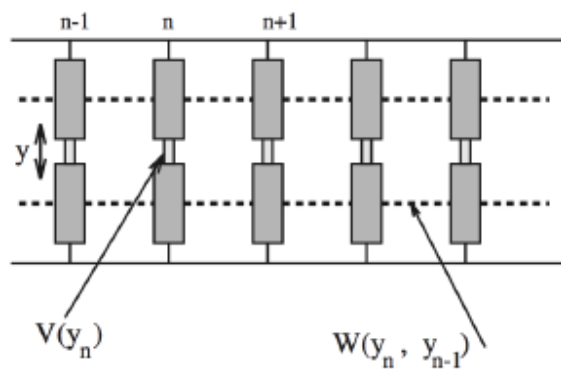


X rays



M. Peyrard, Nonlinearity **17** (2004) R1

$$H_y = \sum_{j=1}^{N-1} W(y_j, y_{j+1}) + \sum_{j=1}^N V_j(y_j)$$



$$W(y_j, y_{j+1}) = \frac{1}{2}k \left[1 + \rho e^{-b(y_j + y_{j+1})} \right] (y_j - y_{j+1})^2$$

$$V_j = D_j \left(1 - e^{-\alpha_j y_j} \right)^2$$

7 parameters:

k = speed of sound

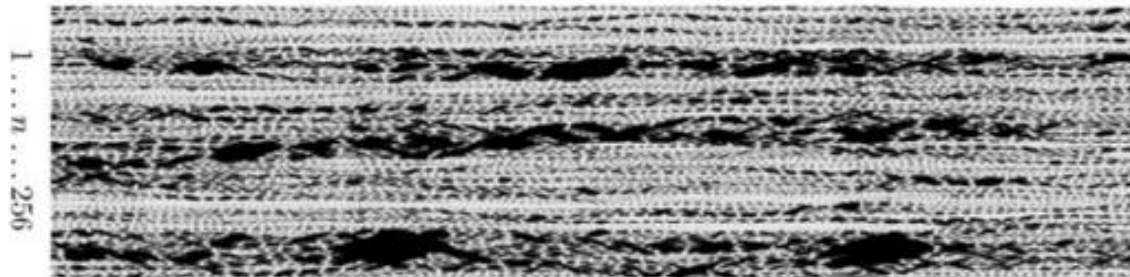
ρ = relative strength of stacking

b = inverse range of stacking

$D_{AT, GC}, \alpha_{AT, GC}$ = Morse potentials

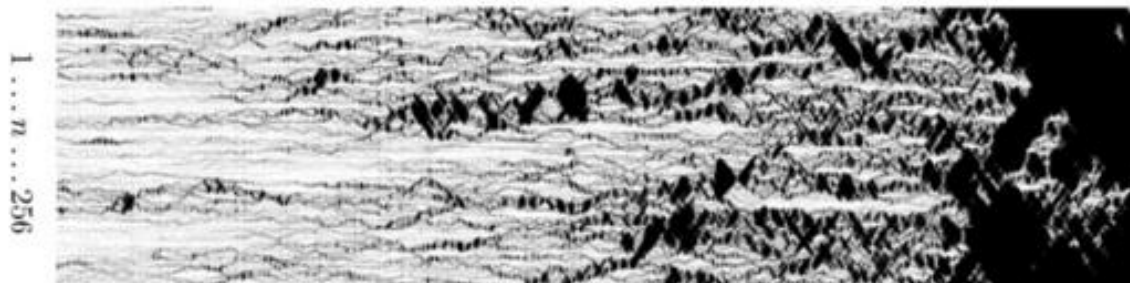
Simulations in M. Peyrard, Nonlinearity **17** (2004) R1

Time

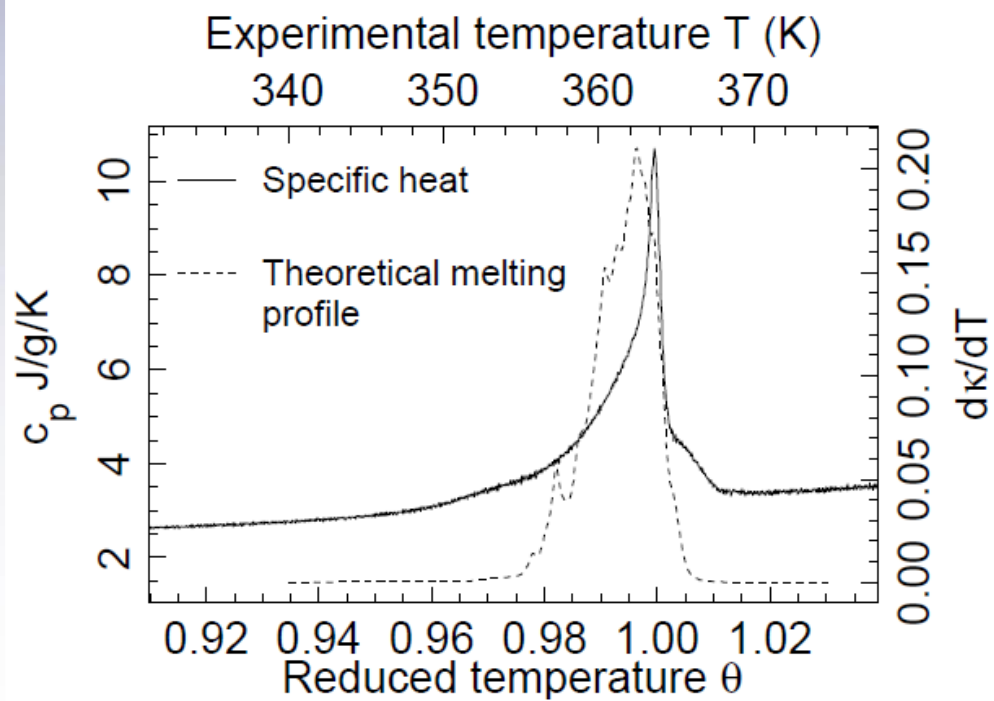


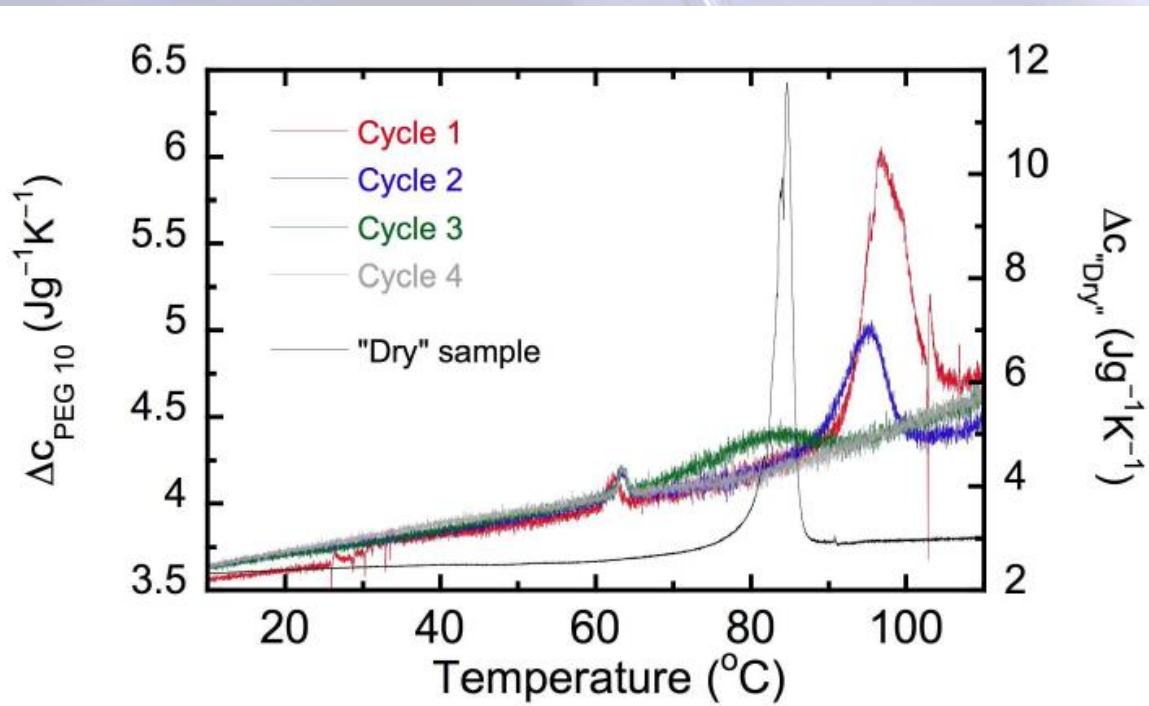
T fixed at 340K

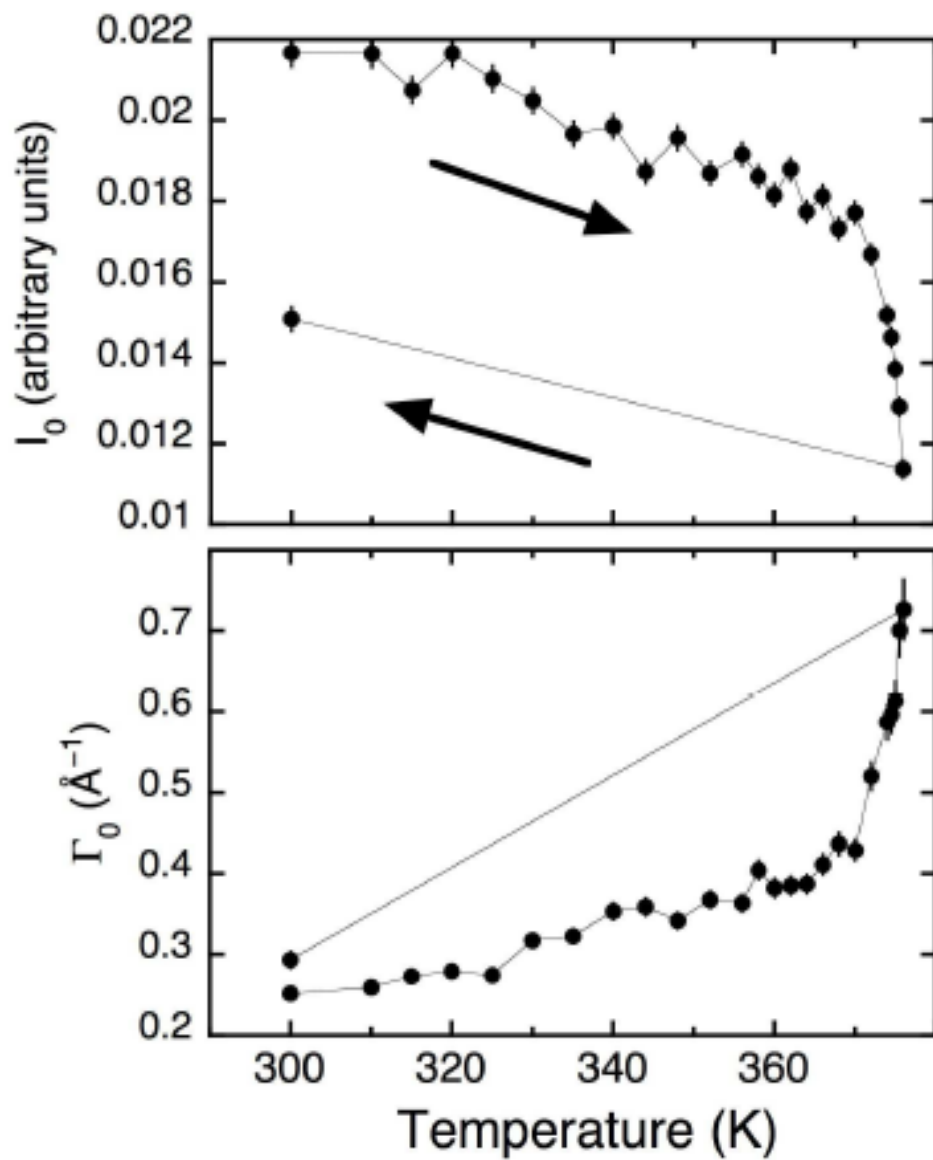
Temperature/Time

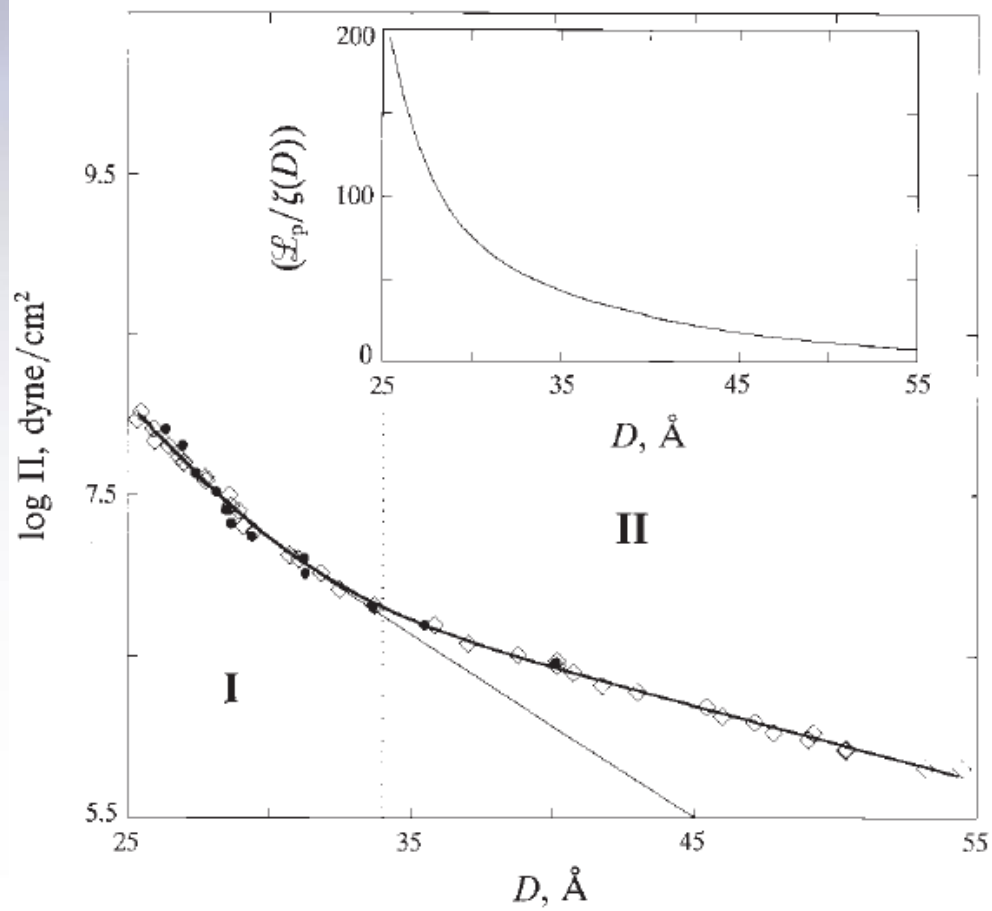


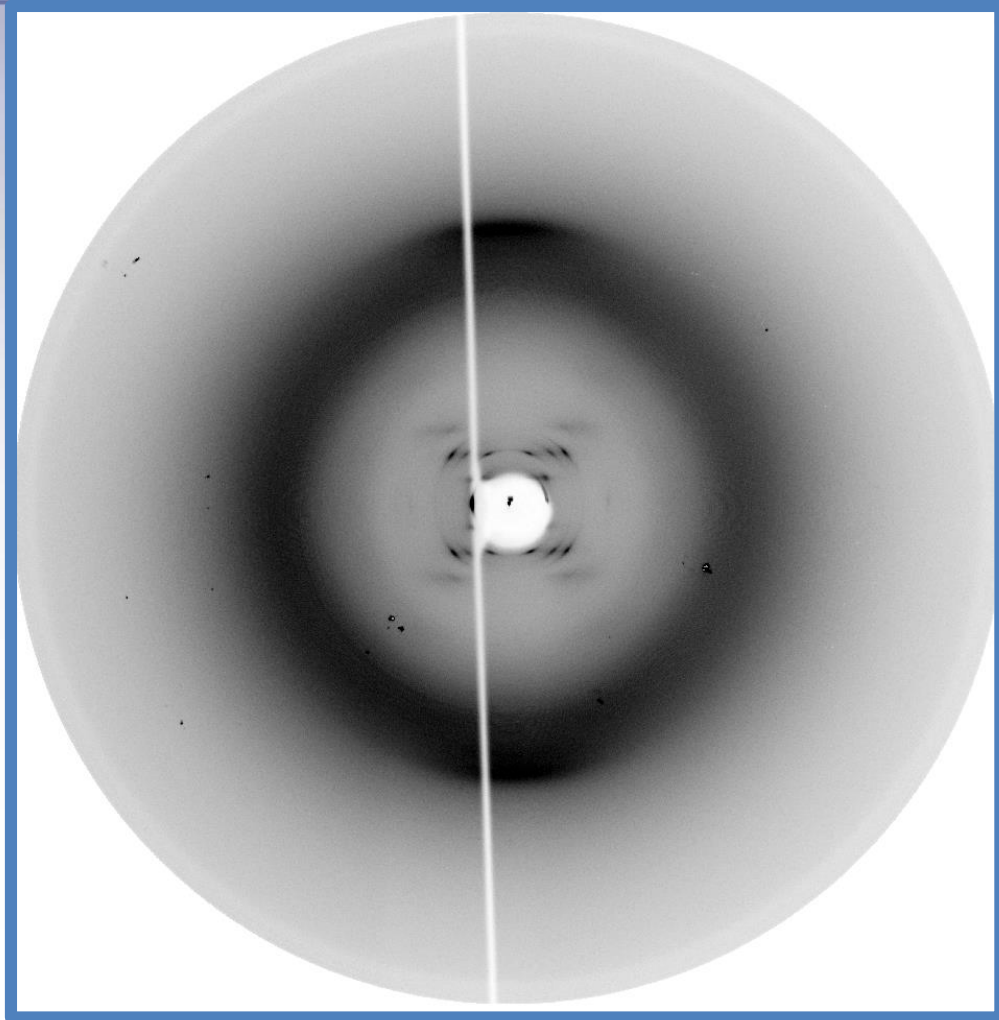
T increasing with time











PEG 6k 17% 12-days

